

NUCLEAR PHYSICS (5)

- a) NATURAL & UNNATURAL SYSTEMS
IN PHYSICS
- b) THE NUCLEAR FORCE &
IT'S PROPERTIES
- CHECK A
FEW
EXAMPLES
- QUANTITATIVE
PROPERTIES

RECAP | [Physical systems \rightarrow Scales]
(scales) $\rightarrow Q = \{Q_1, Q_2, Q_3, \dots\}$
(momentum)

3) Natural systems: one scale dominates

$\bullet \rightarrow Q_1 \ll Q_2, Q_3, Q_4$ $\bullet \rightarrow$ this scale is special

$$\langle \underline{1} \rangle = c_0 \underline{Q_1^d}, \quad \underline{c_0} \sim \underline{O(1)}$$

\rightarrow all observables are natural
in terms of the dominant scale

But we also have ...

2) Unnatural systems: two or more scales
fight with each other

$$\boxed{Q_1 \cup Q_2} \ll Q_3, Q_4, \dots$$

or $\boxed{Q_1 \cup Q_2 \cup Q_3} \ll Q_4, Q_5, \dots$

or any other variation

It's compute
observable

$$\langle \psi | \hat{O} | \psi \rangle = (\text{it's complicated}) \rightarrow \left[\begin{array}{l} \text{there might} \\ \text{be fine-tuning} \end{array} \right]$$

the calculation could
be a surprise

FINE-TUNING

Basically, a cancellation among the effects
of several scales

is { (similar to obtaining a small number
as the difference of two big numbers)

$$A - B = C$$

(but sometimes it
can happen)

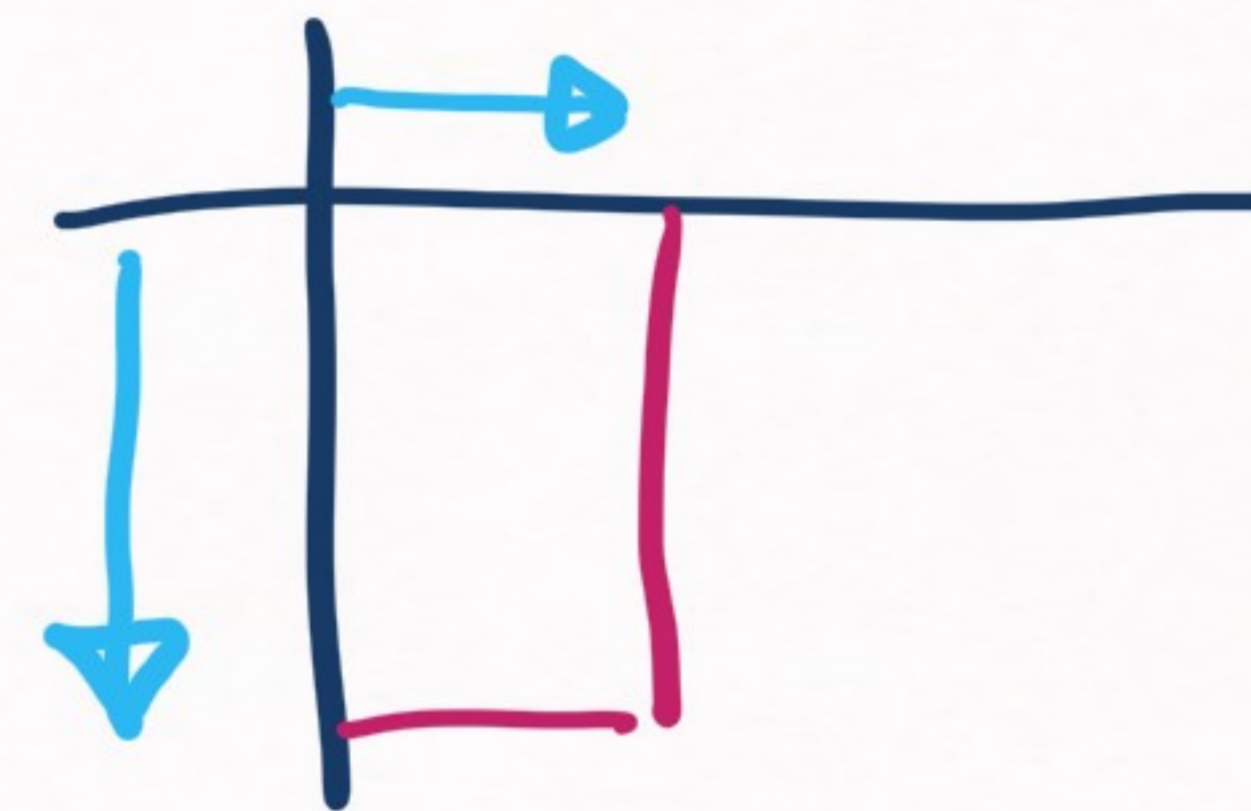
$$\left. \begin{array}{l} A \sim O(10^6) \\ B \sim O(10^6) \\ C \sim O(1) \end{array} \right\}$$

usually this
will not
happen

FINE-TUNING → THE SQUARE WELL

a) Two parameters : range & strength

⇒ Two scales : $\{a, R_s\}$



b) When we have $\boxed{\frac{a}{R_s} = \frac{\pi}{2}}$ there is a zero-energy bound state

→ c) Condition very specific : $\frac{a}{R_s} = \frac{\pi}{2} (1 + x)$

⇒ $x \lesssim 0.2$ for this to happen (fine tuning)

FINE-TUNING \rightarrow [NUCLEAR FORCES]

a) np systems w/ $S = 0, 1$ (spin) (S -waves or $L=0$)

b) For $S=1$, we have the deuteron

$$|E_B| \sim \underline{2 \text{ MeV}}$$

$$\underline{E_B} = \underline{\langle T \rangle} + \underline{\langle V \rangle}$$

$$|\langle T \rangle|, |\langle V \rangle| \sim \underline{50 \text{ MeV}}$$

Fine-tuning of $1/25$ ($2/50$)

c) For $S=0$, we have the virtual state (labeled bound state)

$$|E_V| \sim \underline{0.07 \text{ MeV}}$$

Fine tuning of $\underline{1/700}!!$

[We want to understand
why fine-tuning happens]

↳ TWO ANSWERS :

a) FORTUITOUS / BY CHANCE

↳ just the way nature happens to be

b) "CONSPIRACY" / \exists A MORE PROFOUND REASON

↳ fine-tuning points toward new physics

a) BY CHANCE

FINE-TUNING

→ \exists of eclipses : relative size of the moon and the sun are almost identical on earth

(no fundamental reason for this)

→ Solar system & the location of earth

we just happen to be in the habitable zone of the sun

(no reason, only selection bias)

if this not the case,
we will not be here

[b) "CONSPIRACY"]

FINE-TUNING

→ classical electron radius :

points towards new physics : QM & QFT

→ pion mass : (pion is the lightest hadron)

points towards the ρ meson &
chiral symmetry

→ maybe the Higgs mass

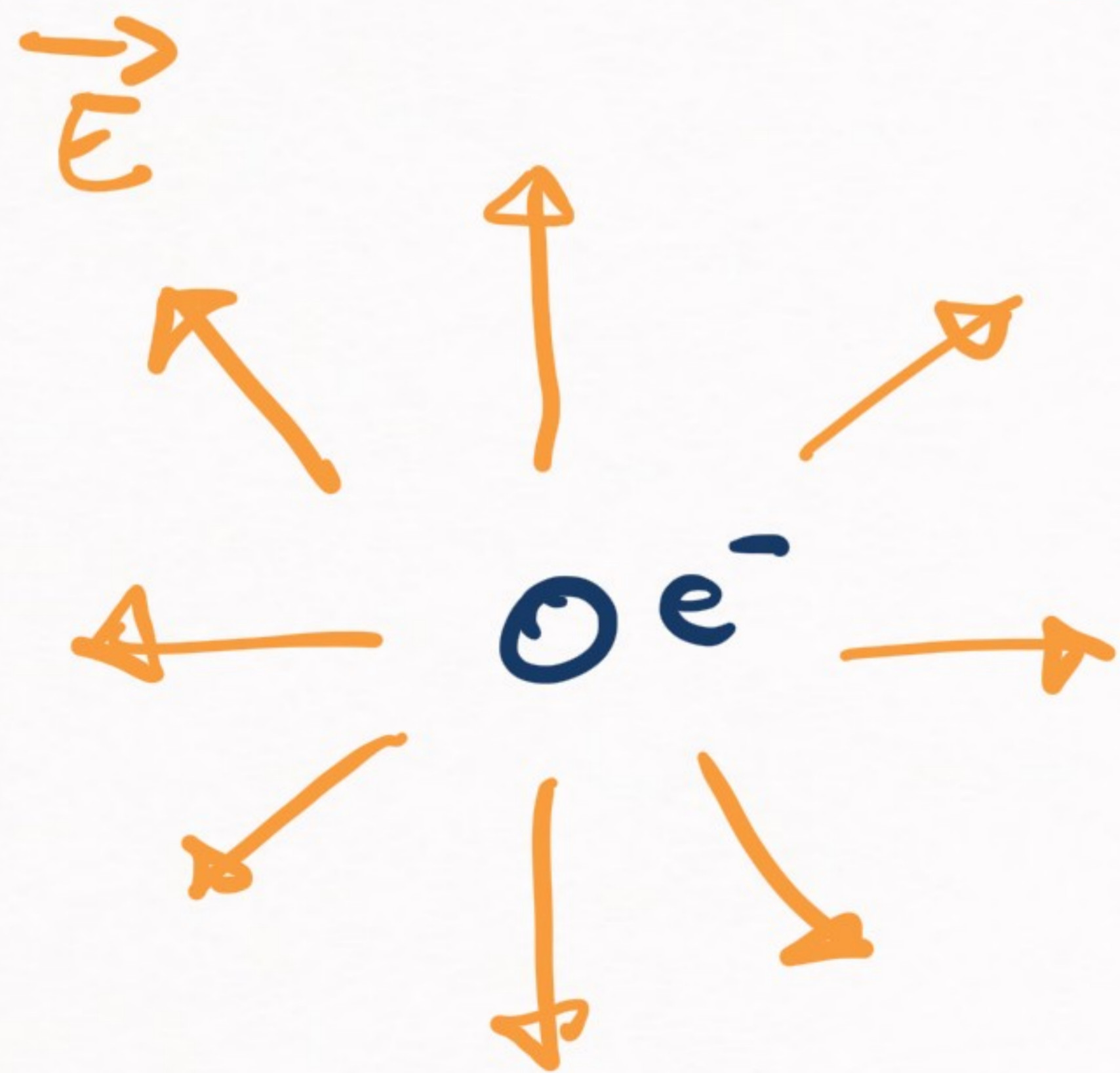
→ maybe the cosmological constant

} (?)

We will review a few examples of
the "conspiracy-type" of fine tuning:

1) CLASSICAL ELECTRON RADIUS

(old calculation in classical
electrodynamics)



Electric field generated by
the electron)

contains energy \rightarrow mass

→ if we take into account the energy
of the electric field around an electron:

$$m(e^-) = \underbrace{m_{\text{bare}}(e^-)}_{\text{"bare" mass}} + \underbrace{m(\vec{E})}_{\text{energy/mass of the } \vec{E} \text{ field}}$$

⇓
mass of the electron if we "switch off"
the electric force

$m(\vec{E}) \rightarrow$ can be calculated

If we assume a point-like electron:

$m(\vec{E}) \rightarrow \infty !!$ $\oplus \rightarrow$ infinite
fine-tuning

$$m(e^-) = \underbrace{m(\vec{E})}_{\infty} + \underbrace{m_{\text{bare}}(e^-)}_{-\infty} = \underline{\text{finite}} \oplus$$

Infinite fine-tuning? → this doesn't make
much sense

What are the alternatives?

Condition: $m_{\text{bare}}(\bar{e}) \geq 0$

(Looks like a sensible
additional condition)

≡

CLASSICAL RADIUS OF THE ELECTRON:

$$m(e^-) = m(\vec{E}) + m_{\text{bare}}(e^-) \rightarrow \boxed{m(e^-) = m(\vec{E})}$$

(≥ 0) ↓


$|\vec{E}| \propto \frac{1}{r^2} \rightarrow$ this energy will depend on the size of the electron

 \rightarrow classical electron radius

$$\boxed{r_{cl}(e^-) = \frac{\alpha}{m_e}}$$

Classical radius of the electron:

$$r_{ce}(e^-) = \frac{\alpha}{m_e} \approx \underline{\underline{2.8 \text{ fm}}}$$

What interpretation
for this number?

\Rightarrow new physics
before we reach r_{ce}

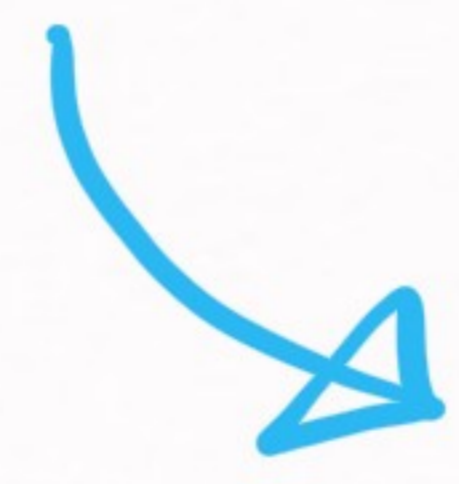
$r > r_{ce}$ \rightarrow something new (not classical)
must happen

Indeed we have new physics for $\boxed{r \geq r_{cl}(e^-)}$

1) Quantum mechanics: $r \sim a_B \sim \frac{1}{m_e \alpha} \gg \frac{\alpha}{m_e}$

2) Quantum field theory: $\left(\frac{1}{\alpha^2} \text{ larger than } r_{cl}(e^-) \right)$

$$r \sim \frac{1}{m_e} \gg \frac{\alpha}{m_e}$$



gives us a new bound
for new
physics

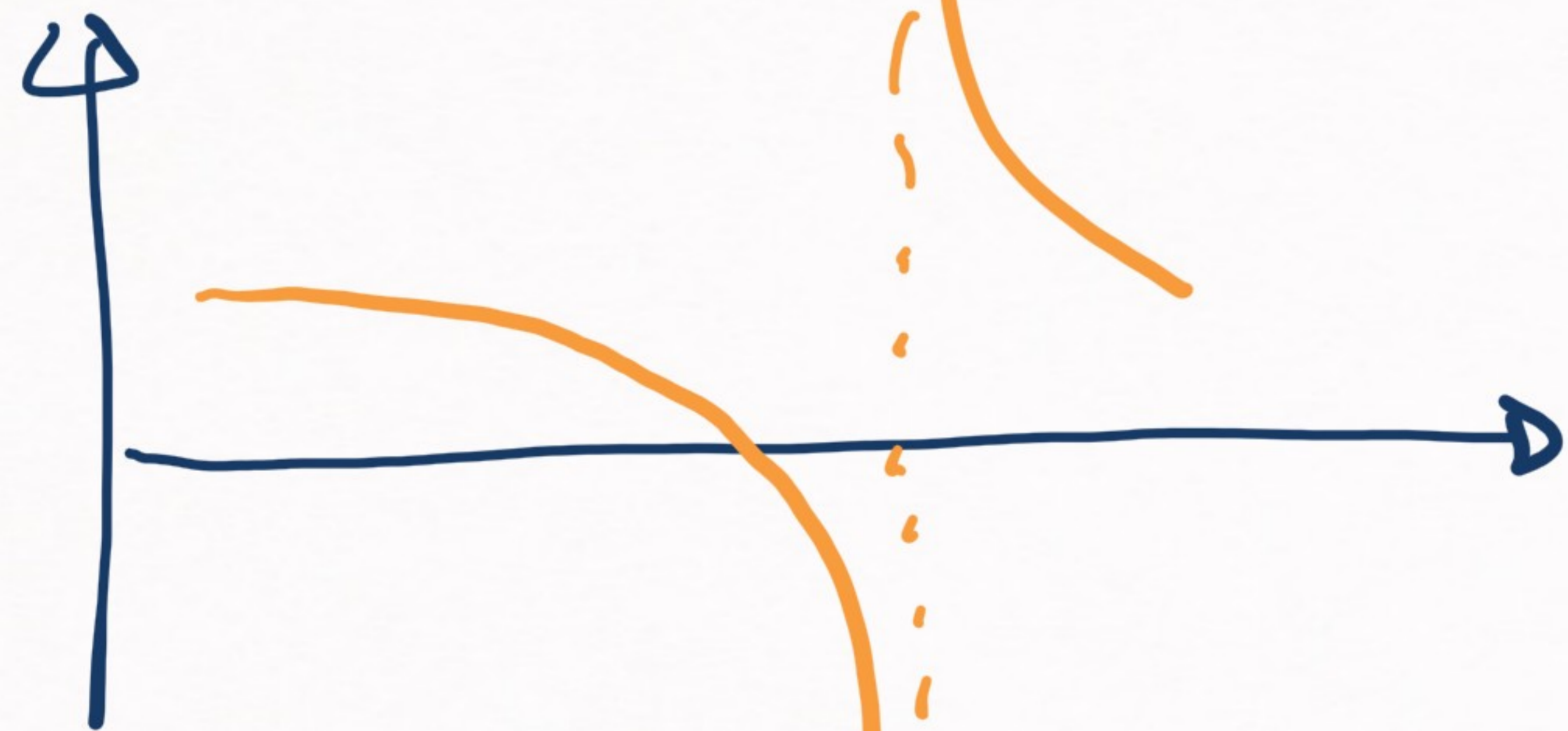
$\left(\frac{1}{\alpha} \text{ times larger than } r_{cl}(e^-) \right)$

$m(e^-)$ in quantum field theory

at which energies do we have $m_{\text{bare}}(\bar{e}) \geq 0$

or the condition $m_{\text{bare}}(\bar{e}) \geq -\infty$

m_{bare}



it's called a cutoff
 Λ (max. energy/momentum
in our calculation)

pole

QFT → At which energies do we have a problem?

LANDAU POLE

(renormalization of electron charge goes to infinity)

$$\Lambda \sim 10^{280} \text{ MeV}$$

$$r \sim \frac{\hbar c}{\Lambda} \sim \frac{200}{10^{280}} \text{ fm} \\ \sim 2 \cdot 10^{-278} \text{ fm}$$

LANDAU POLE

$$\left. \begin{array}{l} r \sim 2 \cdot 10^{-27} \text{ fm} \\ \Lambda \sim 10^{280} \text{ MeV} \end{array} \right\} \rightarrow \oplus$$

⊙ → Again, we should expect new physics

$$\text{for } r > 2 \cdot 10^{-27} \text{ fm}$$

(kind of obvious → $\Lambda_{\text{Landau}} \gg \underline{\underline{M_{\text{black}}}}$)

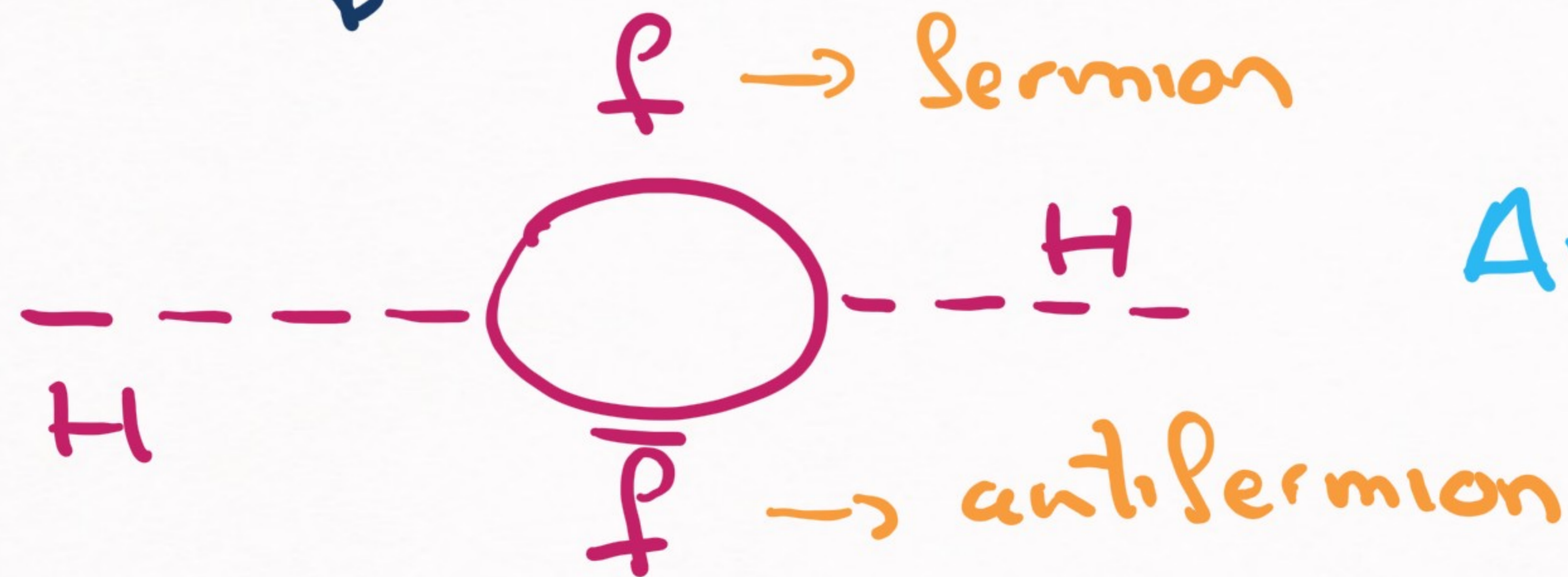
[Quantum gravity]

2) MASS OF THE HIGGS

(ANALOGOUS TO THE PION MASS PROBLEM)

$$m_H^2 = m_{\text{bare}}^2 + \underbrace{\Delta m_H^2}_{\text{like before w/ the electron}}$$

in QFT this comes from loop corrections



$$\Delta m_H^2 \propto -\frac{|\lambda_f|^2}{8\pi^2} [\Lambda^2 + \dots]$$

The point is that corrections to the mass
of the Higgs goes as Λ (the cutoff
of the theory)

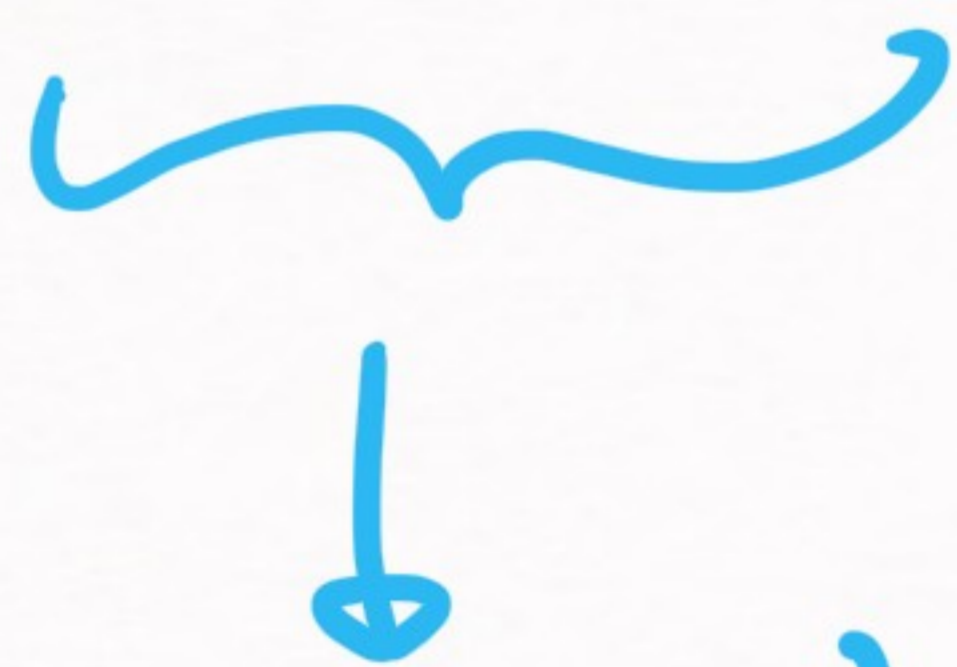
↙ characteristic of spin-0 fields
(Higgs, pion, etc.)

In contrast, fermions receive mass corrections
that go as $\log(\frac{\Lambda}{Q}) \rightarrow$ Landau
pole

Fermions $\rightarrow \Delta m \propto \log\left(\frac{\Lambda}{\bar{a}}\right) \rightarrow$ EXAMPLE:

Landau pole
at incredibly high
energies

Spin-0 particles $\rightarrow \Delta m \propto \Lambda$



super sensitive
to short-range

(grows really fast



problems will appear before)
physics

What is the amount of fine-tuning
for the Higgs?

→ $\frac{m_H}{\Lambda}$ (ratio of the Higgs mass over
the ultraviolet scale
of the standard model)

→ ultraviolet → happens for Λ large (short
-distances)

→ infrared → happens for Λ small
(long-distances)

How large is Λ ?

$$\frac{m_H}{\Lambda}$$

a) A possibility is that the standard model is valid up to the Planck scale (distances at which we can see quantum gravity)

$$G \sim \frac{1}{M_{\text{Planck}}^2} \rightarrow M_{\text{Planck}} \simeq 10^{18} \text{ GeV}$$

Newton's constant (gravity)

$$\frac{m_H}{\Lambda} \sim 10^{-16}$$

b) Second possibility \rightarrow standard model stops to be valid sooner

$$\frac{m_H}{\Lambda} \sim \frac{1}{10} \text{ or } \frac{1}{100} \text{ or } \frac{1}{1000} \rightarrow \text{depends on } \Lambda$$

EXAMPLE \rightarrow SUSY (supersymmetry) $\Lambda \sim M_{\text{SUSY}}$

Usual estimation: $M_{\text{SUSY}} \sim (1-10) \text{ TeV}$

$$\frac{m_H}{\Lambda} \sim \left(\frac{1}{10} - \frac{1}{100} \right) \rightarrow \text{less fine-tuned}$$

PROBLEM | →

We don't know what new physics lies beyond the standard model

$$\frac{m_H}{\Lambda} \gg 10^{-16}$$

↳ ∃ infinite reasons why this is not so fine-tuning

BOTTOM-LINE \rightarrow $\frac{m_H}{\Lambda}$ probably indicates
new physics

$\left. \begin{array}{l} \text{ } \\ \text{ } \end{array} \right\} (\Lambda \sim 10 m_H, 100 m_H, 1000 m_H$
 $\text{etc.})$

any of these options look better than $10^{16} m_H$
 \sim

3) COSMOLOGICAL CONSTANT (super fine-tuning example)

Reminder:

$$R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} + \underbrace{\Lambda_c g_{\mu\nu}}_{\text{cosmological term}} = 8\pi G T_{\mu\nu}$$

$[\Lambda_c] = [E^4]$ → What is its expected natural size?

$$\Lambda_c = Q^4$$

with Q a natural scale
in gravity,

one answer:

$$G \sim \frac{1}{M_{\text{Planck}}^2} \rightarrow Q \sim \underline{M_{\text{Planck}}}$$

$$\begin{aligned} \Lambda_c &\sim (10^{18} \text{ GeV})^4 = (10^{27} \text{ eV})^4 \\ &= 10^{108} \text{ eV}^4 \end{aligned}$$

a) natural estimation of Λ_c

$$\Lambda_c \sim 10^{103} \text{ eV}^4$$

b) but we can measure Λ_c from
the acceleration of the expansion
of the universe

$$\Lambda_c \sim 10^{-12} \text{ eV} \quad \left| \quad \text{HOLY HOLY !!} \right.$$

The natural estimation fails by a factor of:

$$\left[\frac{\Lambda_c}{M_{\text{Planck}}^4} \sim \frac{1}{10^{120}} \right] \rightarrow \text{fine-tuning of } \Lambda \text{ part}$$

in 10^{120} !!!

Looks improbable

→ Requires an explanation

Explanations:

- 1) **Observation bias**: we will only be alive in universes w/ certain values of Λ_c
- 2) **Multiverse**: more than 10^{120} universes it's likely that there will be fine-tuned universes like ours
- 3) **New physics**: QFT understanding of vacuum energy

→ We don't know the answer &

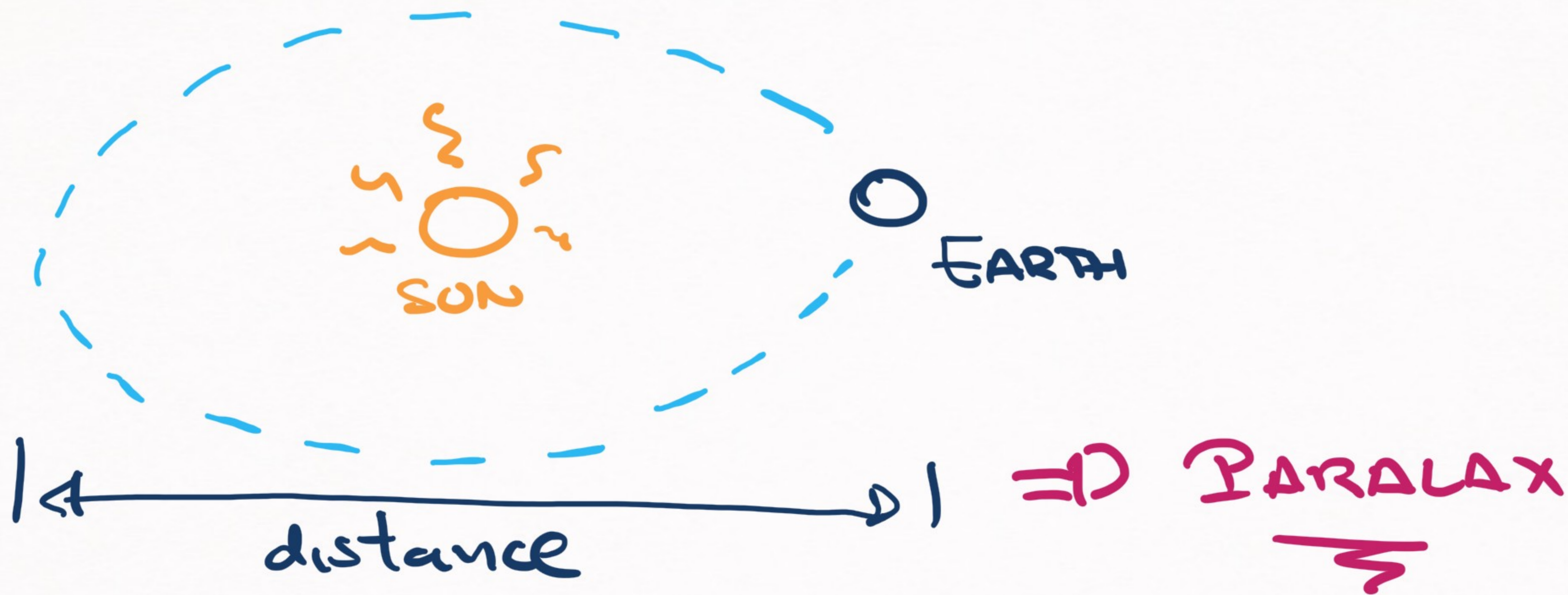
Sometimes assuming quantities must be natural
leads us to wrong answers



EXAMPLE: TYCHO BRAHE'S MODEL
OF THE SOLAR SYSTEM

KEPLER → HELIOCENTRIC MODEL

BRAHE → HAD THE GOOD OBSERVATIONS





PARALLAX: small change
in the relative position
of this star
↯



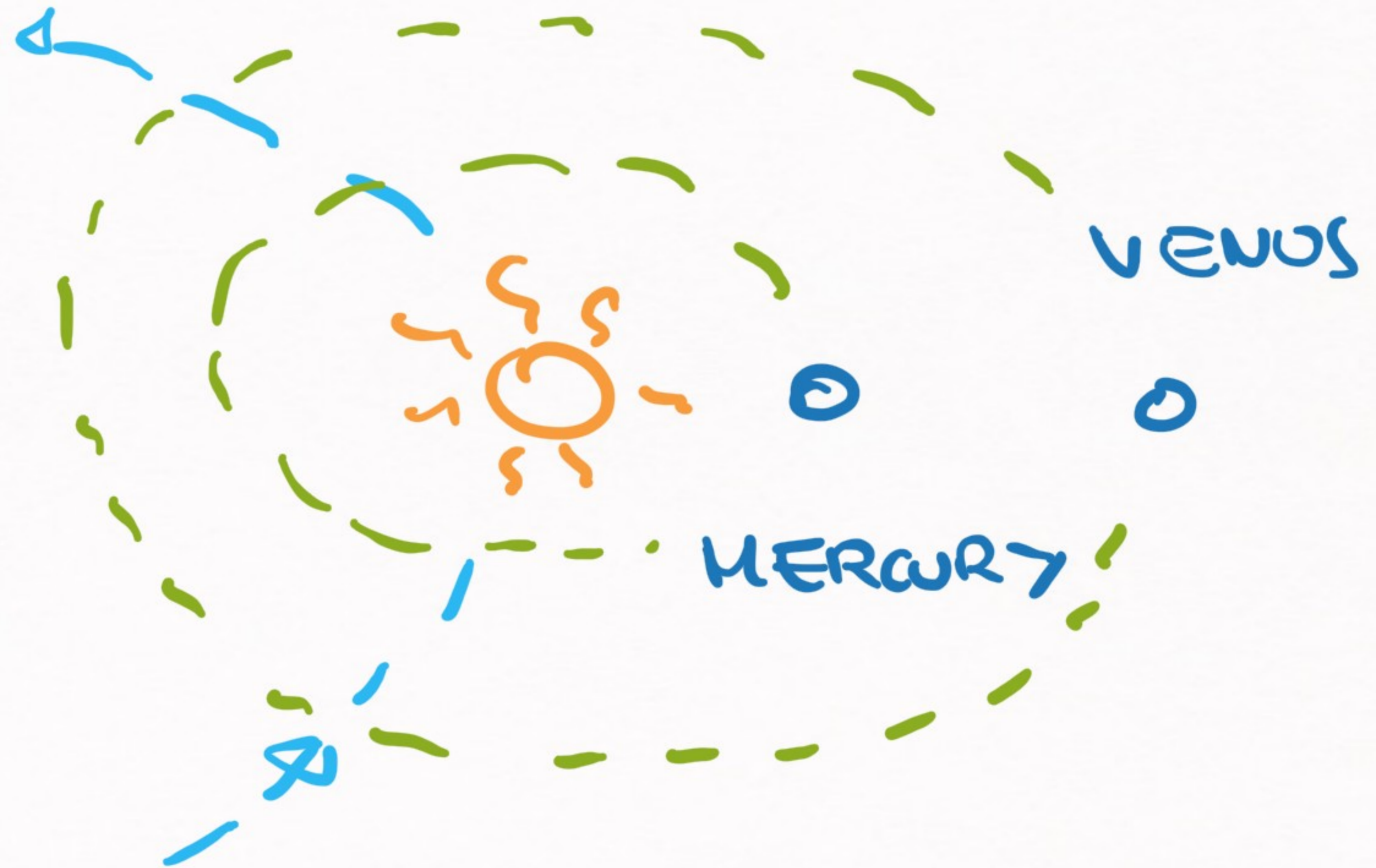
BRADY → DIDN'T OBSERVE ANY PARALLAX

a) Stars really far away

b) Earth is at the center of the solar system { → ⊙

• → choose Ptolemy for this one because otherwise the distances involved would be too unnatural

BRADHE MODEL :



WEIRD MODEL → BASIS WAS NATURALNESS

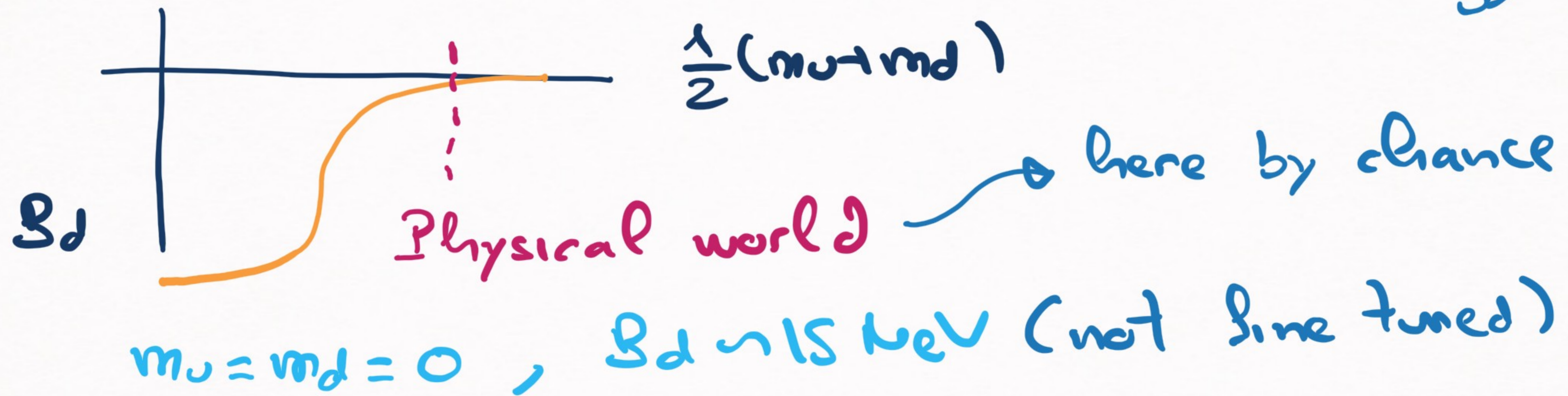
(TAKE HOME MESSAGE)



- FINE-TUNING, WHY?
- 3) Sometimes this requires an explanation
 - 2) Sometimes this is just chance & there is no further explanation
~

FINE-TUNING IN NUCLEAR PHYSICS:

- a) binding energy of deuteron \rightarrow too small
 \rightarrow by chance (probably)
 $\rightarrow m_u, m_d$ quark masses \rightarrow coupling w/
the Higgs



b) mass of the pion

$$m_u, m_d \ll m_s, m_c, m_b, m_t$$

(these two quark masses are small)

$$m_u = m_d = 0 \rightarrow \boxed{m_\pi = 0}$$

\Downarrow
 \exists symmetry, for this
 \searrow (not chance)

→ enough for now about fine-tuning

PROPERTIES OF
THE NUCLEAR FORCES

Nuclear forces \rightarrow forces among neutrons & protons

nucleon \rightarrow neutron, proton



PROPERTIES

- 1) Nuclear forces have a short range
- 2) At "intermediate" distances, nuclear forces are attractive $(1-2) \text{ fm}$
 \searrow

3) At short distances the nuclear force
is repulsive $\leq 0.7 \text{ fm}$

4) Nuclear forces do not distinguish
neutrons & protons (\rightarrow nucleons)

5) Nuclear forces are not central

$$V(\vec{r}) \neq V(|\vec{r}|)$$

(dependence on direction)

PROPERTIES (IN MORE DETAIL)

§) SHORT-RANGED (\neq SHORT RANGE)



I will explain the difference

two types of potential (check previous lessons)

1.a) LONG-RANGED: Coulomb, van der Waals

$V(r) \propto 1/r^n \rightarrow$ extends to infinity

1.b) SHORT-RANGED: Yukawa, square well

$$V(r) \propto \frac{e^{-mr}}{r^n}, \quad V(r) \sim \frac{\Theta(a-r)}{r^n}$$

SHORT-RANGED FORCES \rightarrow Easier to handle
 \searrow

Difference between short-ranged & short range
long-ranged & long range

SCALES:

short & long range

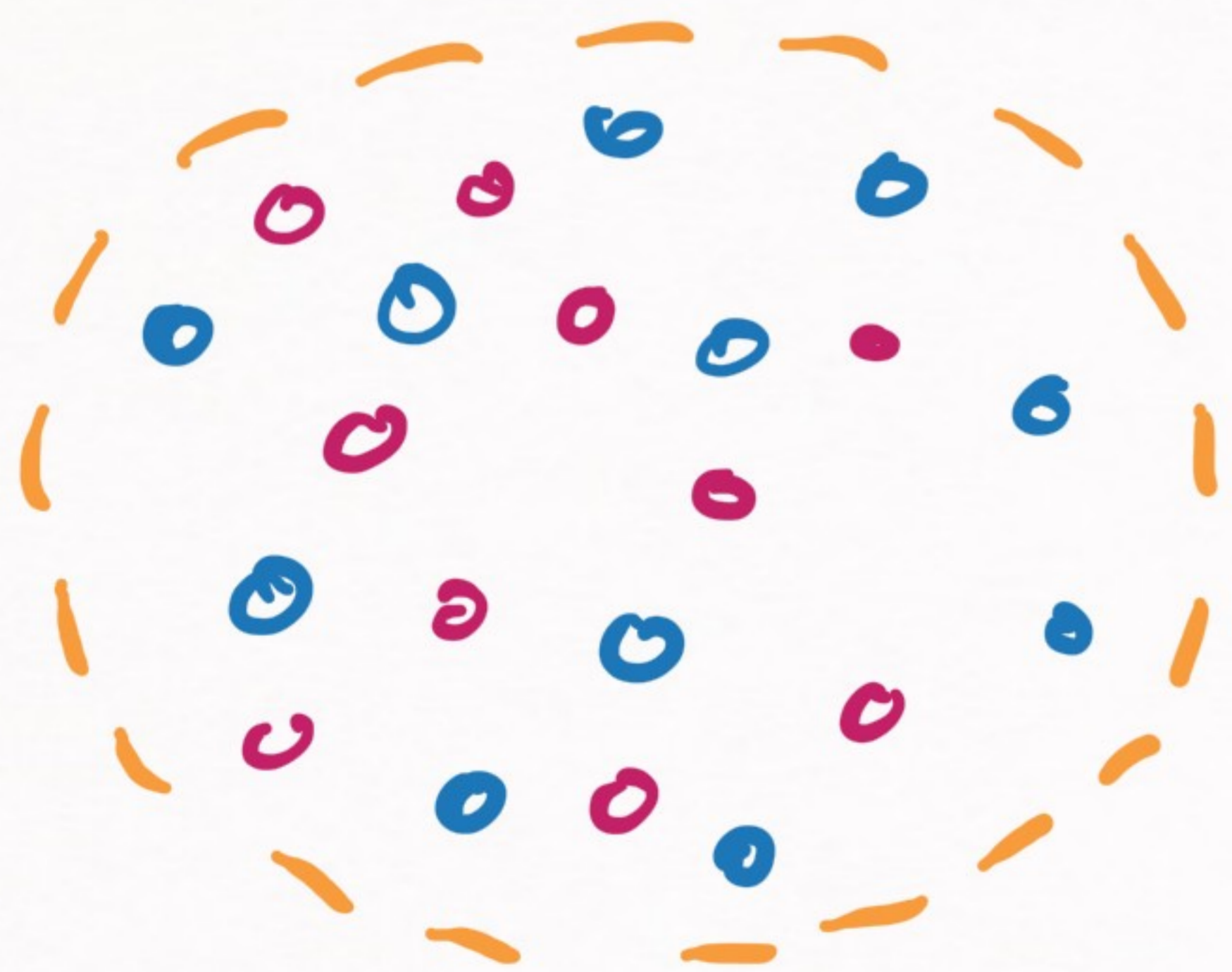
(pedantic, do not assume that other people will differentiate them)

$$V(r) = \underbrace{V_S(r)}_{R_S < R_L} + \underbrace{V_L(r)}$$

\searrow
Only language, not physics)

How DO WE KNOW THAT NUCLEAR FORCES
ARE SHORT-RANGED?

→ Binding energy of nuclei:



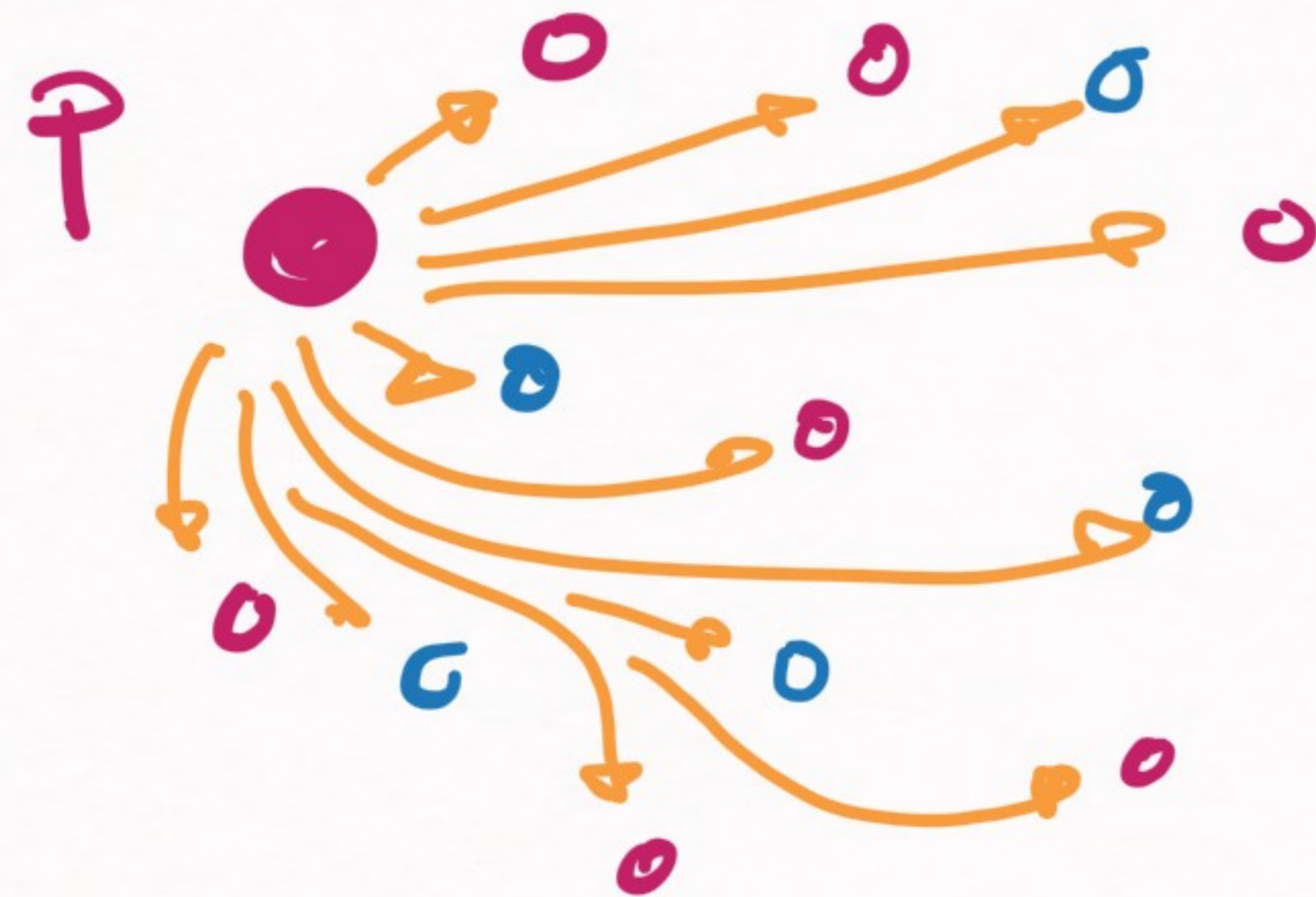
$B = B(A)$ → depends on
range of forces
of nucleons

$$B = B(A)$$

A is number of nucleons

a) if you have a force w/ infinite range (e.g. gravity)

$$B \propto \frac{A(A-1)}{2} \quad (B \propto A^2)$$



Every particle interacts with all others

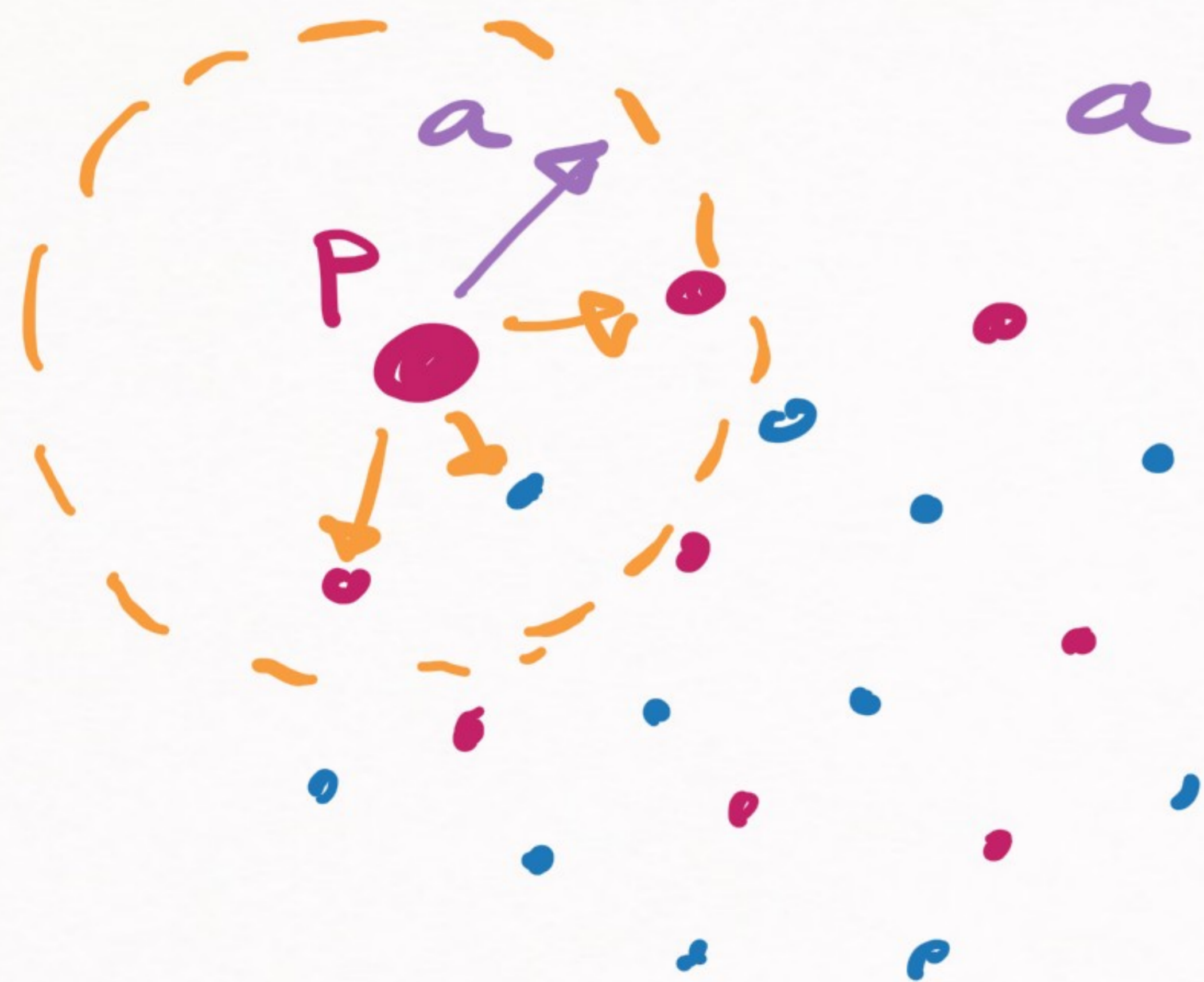


$$B = B(A)$$

A is number of nucleons

b) if we have a short-range force:

$$B \propto A \rightarrow \text{why?}$$



$a \rightarrow$ range of the force

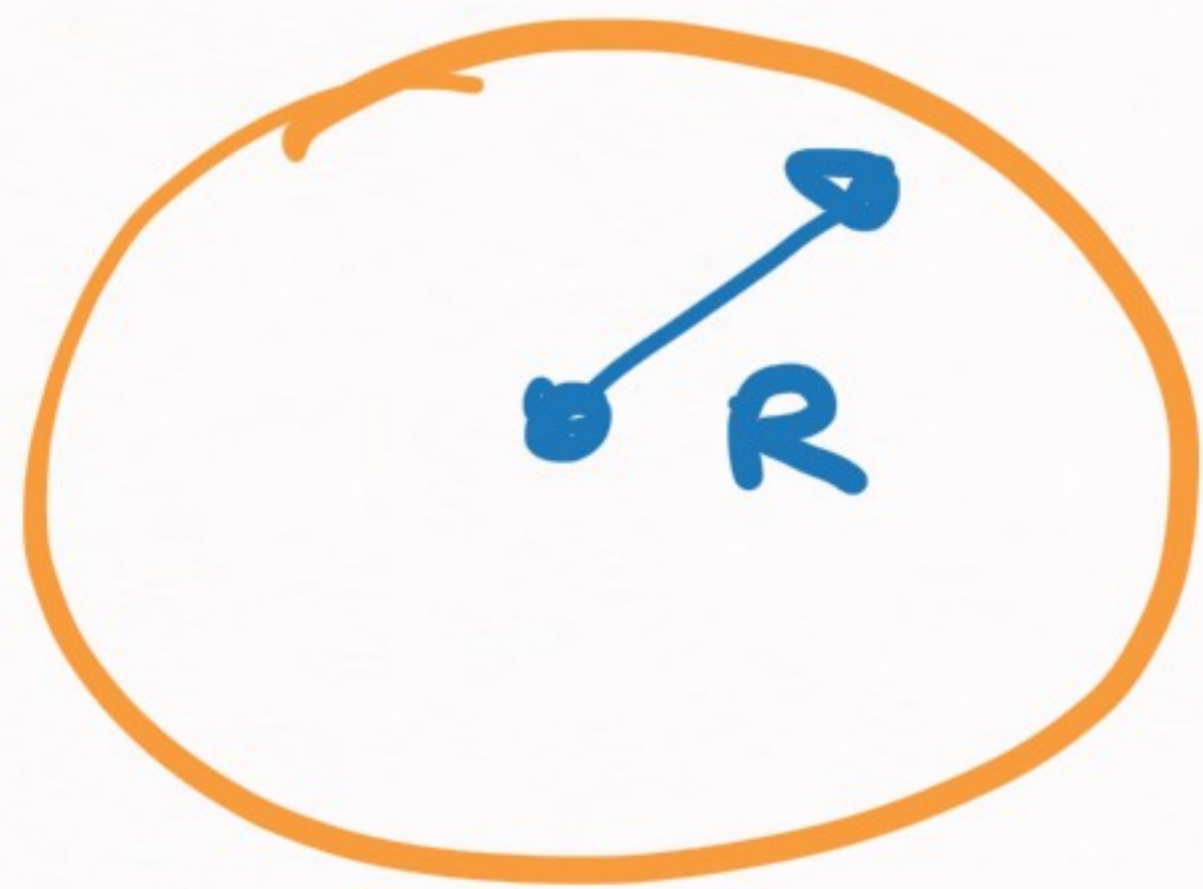
(outside the range, the interaction is practically zero)

Each particle only interacts with the closer neighbors

$$\underline{B \propto A}$$

Exercise to understand this idea:

TO CALCULATE THE GRAVITATIONAL
SELF ENERGY OF A SPHERICAL
MASS DISTRIBUTION



constant
density

$$U \propto M^2$$

HOWEVER, IF WE
MODIFY GRAVITY
WITH A RANGE:

like
nuclear forces

A hand-drawn blue circle representing a sphere. Inside the circle, the text $V \sim \frac{e^{-\lambda r}}{r}$ is written in red. A blue arrow points from the text "like nuclear forces" to the circle.

$\Rightarrow U \propto M$

How DO WE KNOW THAT NUCLEAR FORCES
HAVE A SHORT RANGE?

ANSWER: SATURATION

$$\frac{B}{A} \sim (B - a) \text{ MeV} \quad \left. \vphantom{\frac{B}{A}} \right\} \text{Very good approx. to binding energy}$$

only possible if finite range of interaction

PROPERTIES 1

2) NUCLEAR FORCES ATTRACTIVE AT MEDIUM DISTANCES

(1-2) fm
more or
less
↘

↓
How do we know this?

→ Density of nuclei is more or less constant

$$\rho \neq \rho(A)$$

$$\rho \approx \rho_0$$

PROPERTIES 1

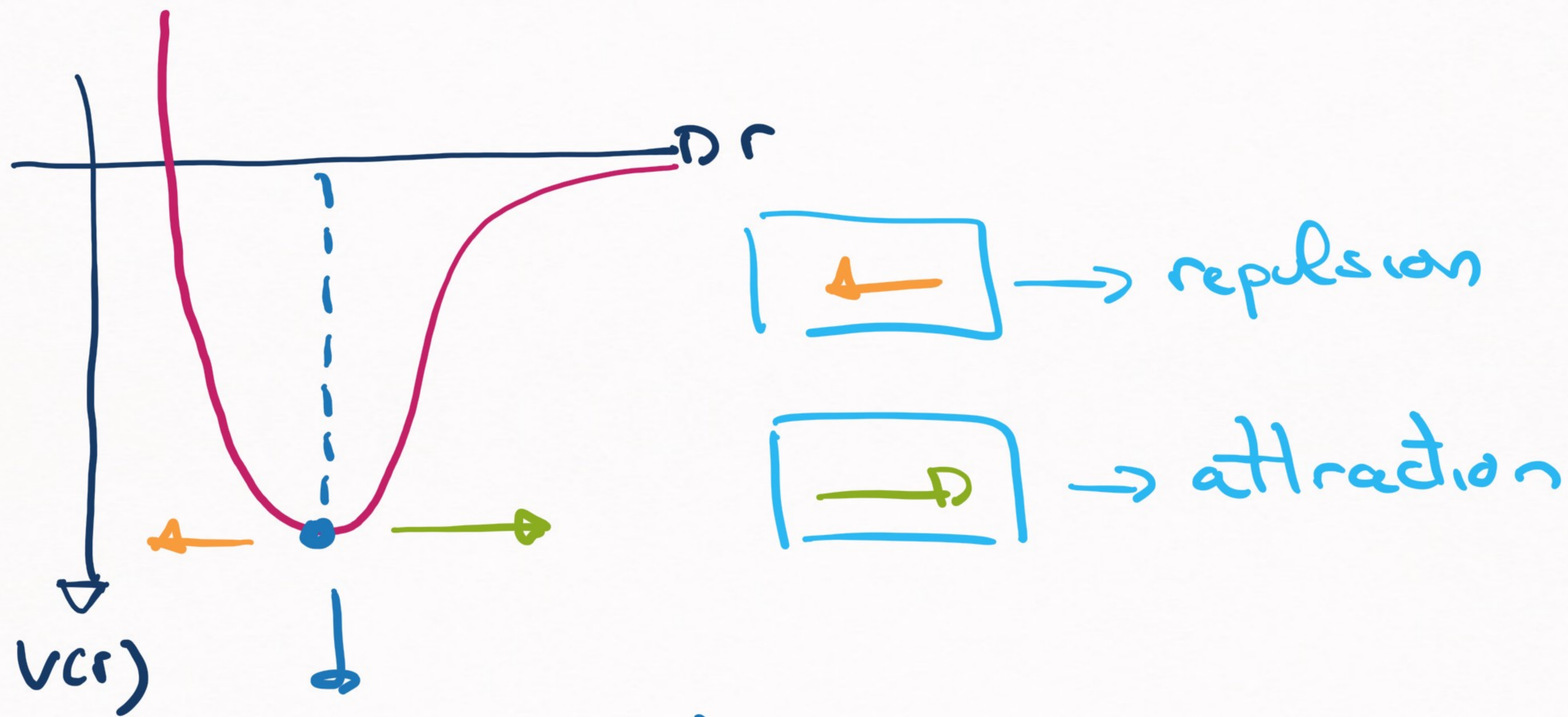
3) NUCLEAR FORCES ARE REPULSIVE
AT SHORT DISTANCES



Again, an important reason is density

$$\rho \approx \rho_0, \quad \rho \neq \rho(A)$$

→ Let's explain 2) & 3) together

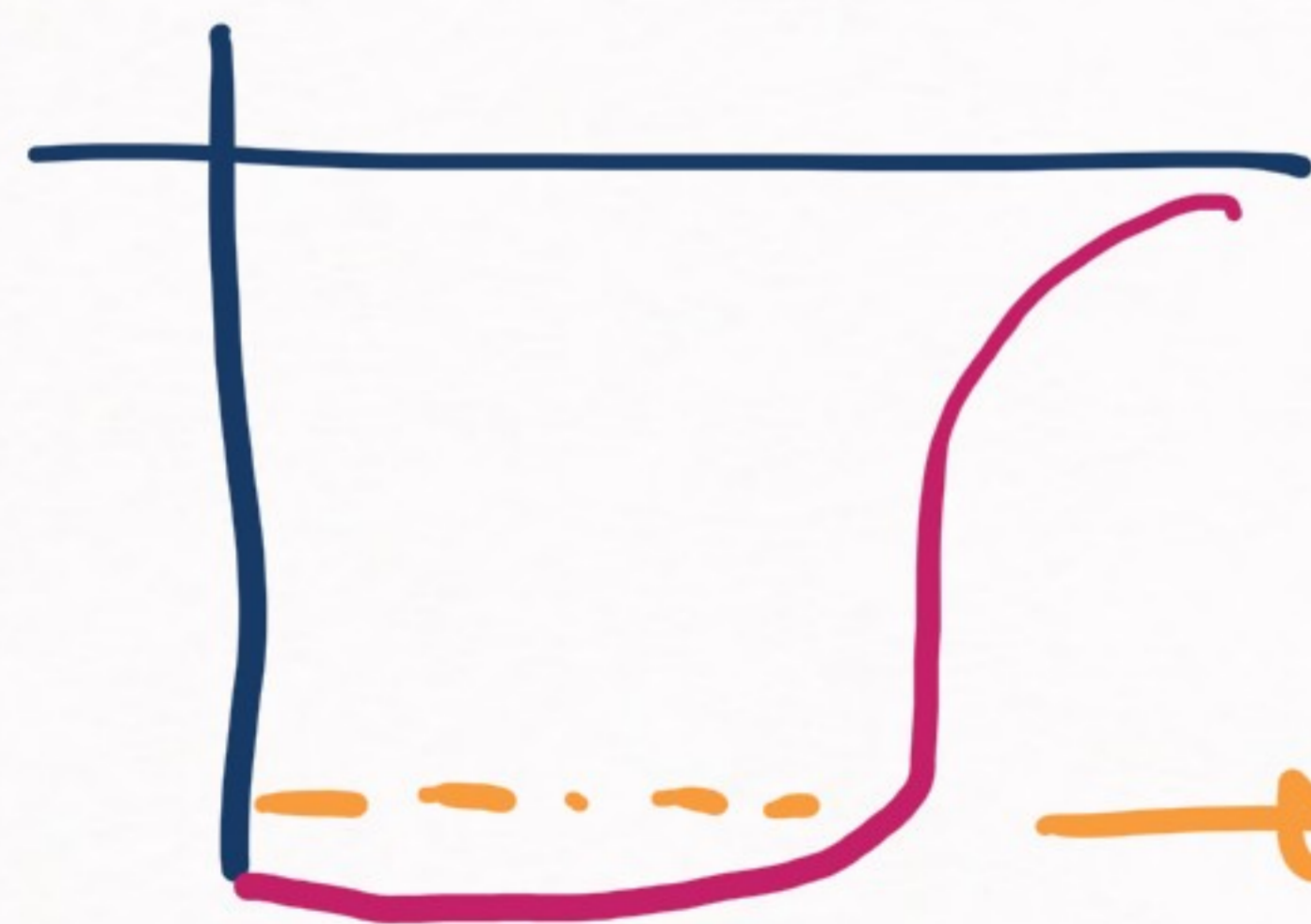


preferred relative
distance between two nucleons

You need a minimum for $r \approx r_0$

(because you want nucleons to be equally spaced)

NOTICE THAT $\frac{dV}{dr} \approx \text{constant}$ DOES NOT IMPLY Δ MINIMUM



→ minimum energy independent of distance

→ We will extend on this
on next lesson

FRIDAY AT 15:50

See you then!