

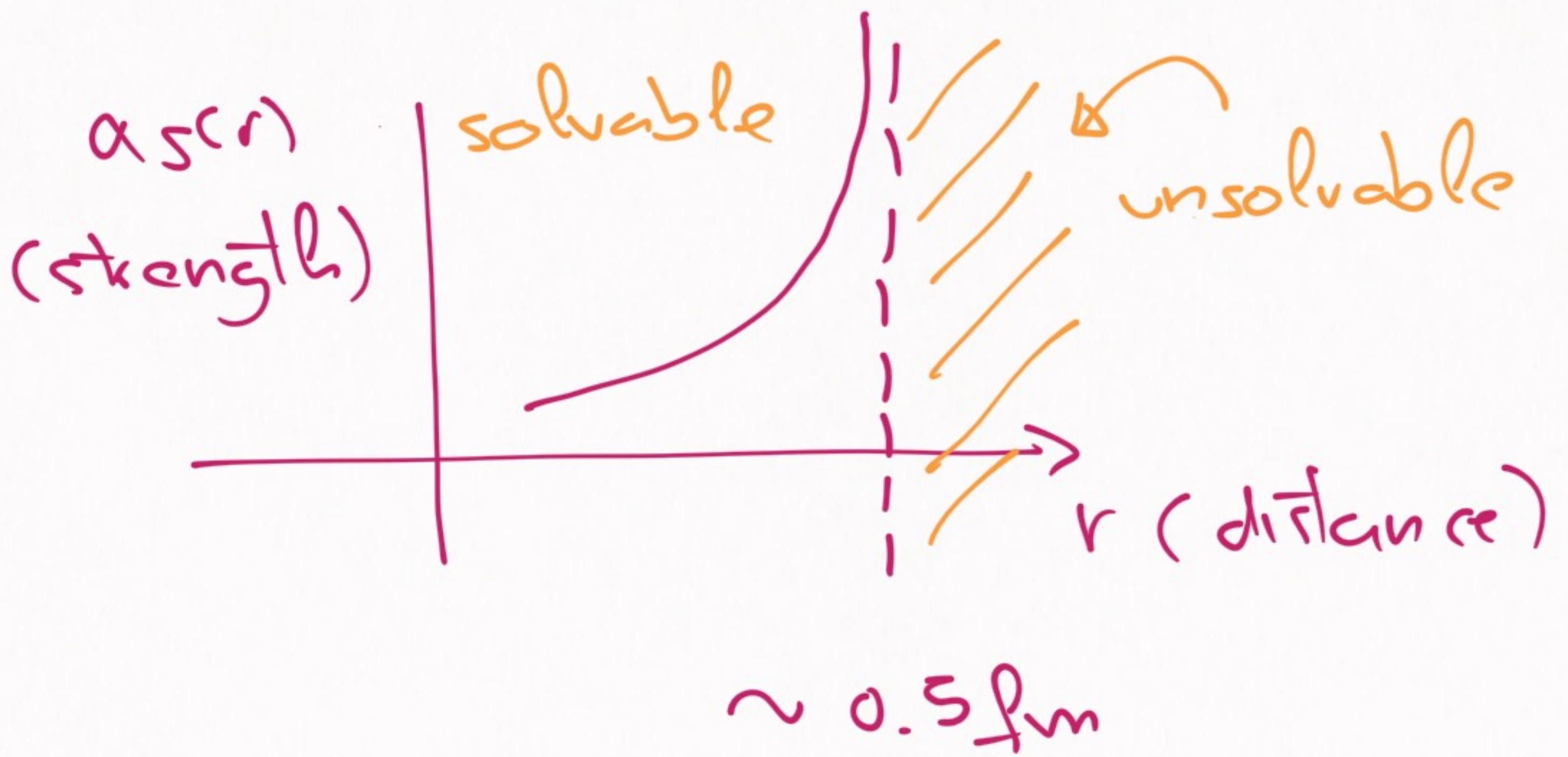
Nuclear Physics (7)



QCD can't be solved
analytically

What are the alternatives?

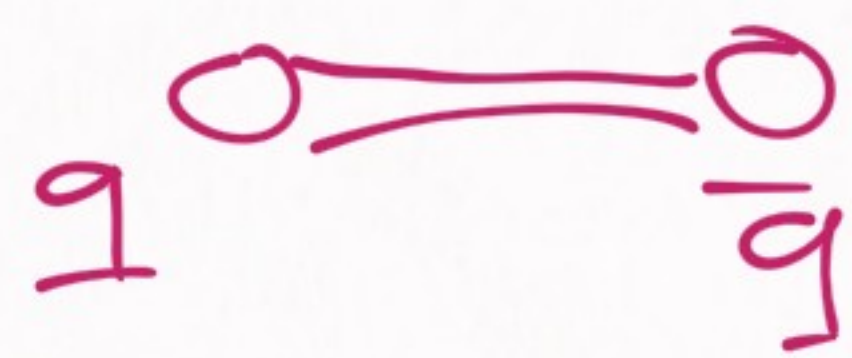
RECAP: QCD has asymptotic freedom & confinement

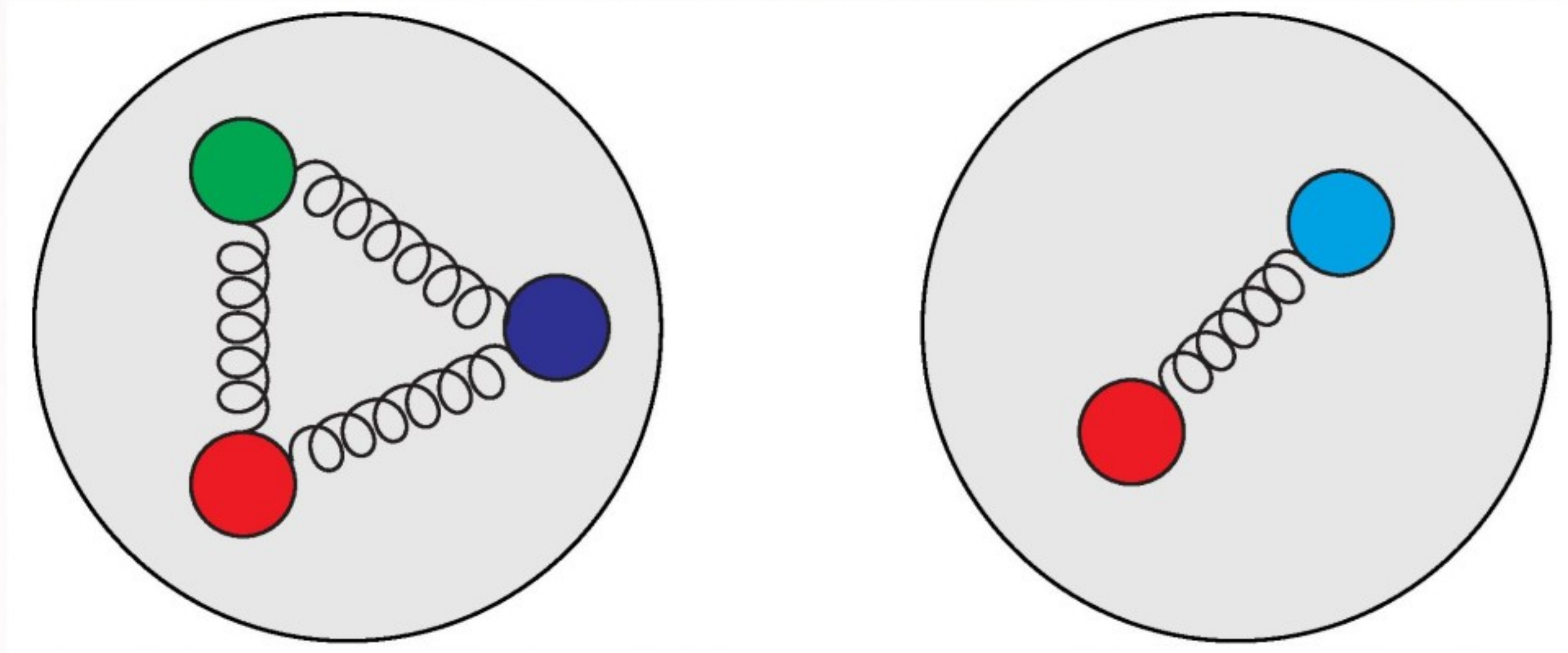


confinement \rightarrow no free quarks



this happens when you try to separate quarks





Baryon

↓
Proton

$r_p \sim 0.85 \text{ fm}$

Meson

↓
Pion

$r_\pi \sim 0.65 \text{ fm}$



Too large for QCD



Alternatives?



Two possibilities are :

1) Lattice QCD

Use supercomputer
to solve QCD



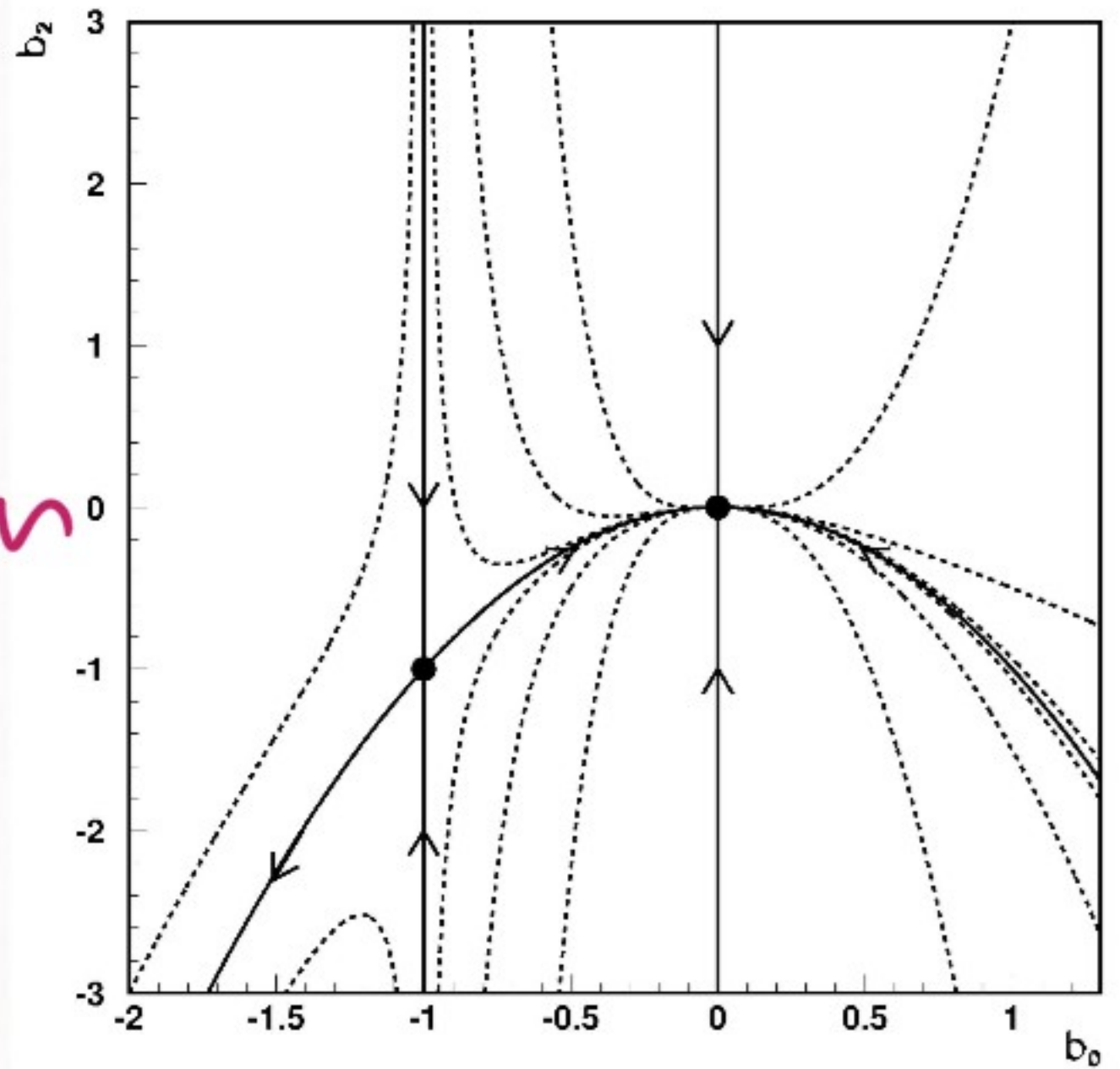
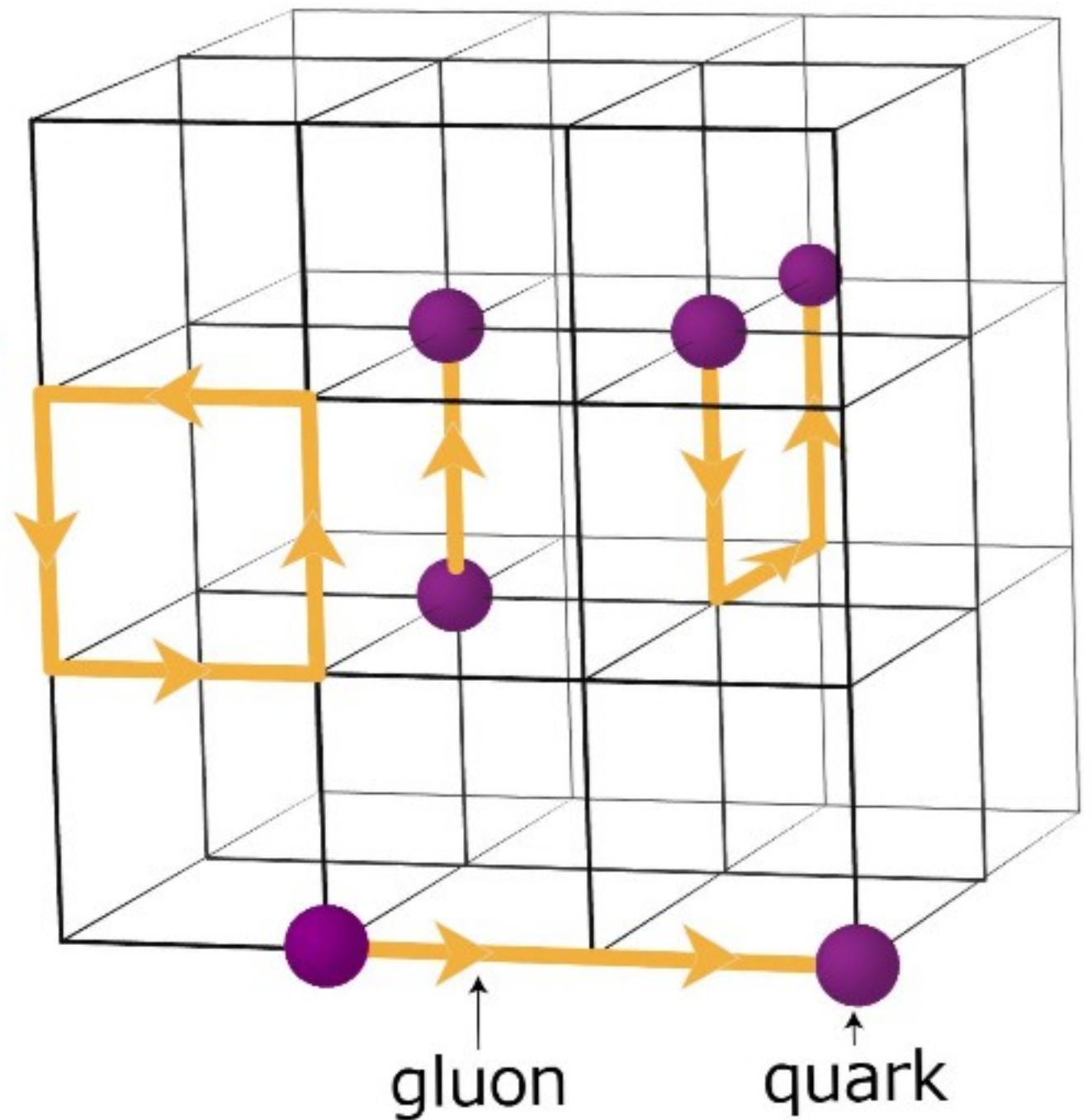
I don't have
a supercomputer



2) Effective Field
Theory

Use renormalization
group analysis
to indirectly
solve QCD

↳ My choice: I only have pen &
paper & a tablet



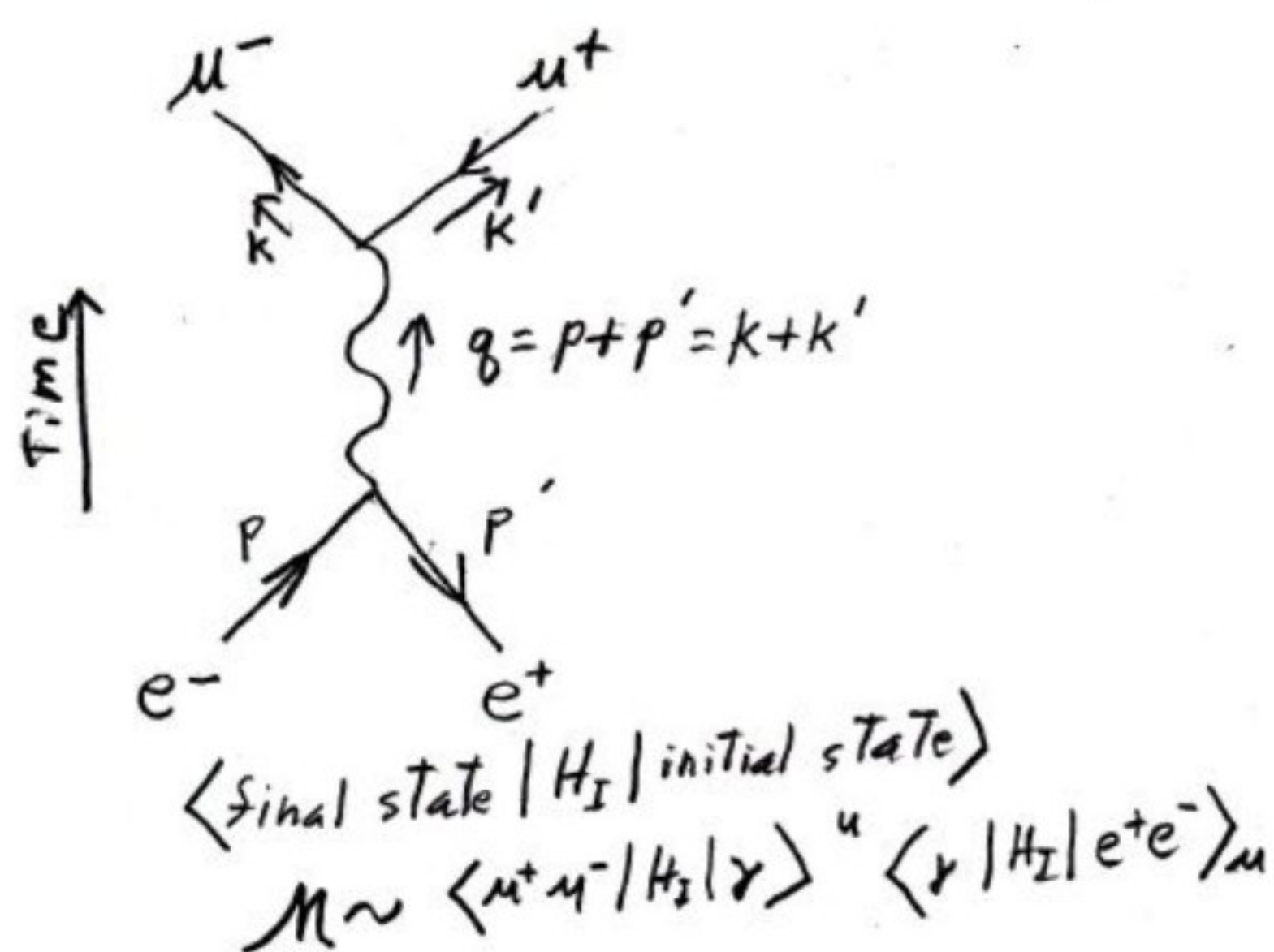
EFFECTIVE FIELD THEORIES (EFTs)

What are they?

→ Let's begin w/ RENORMALIZATION

1) ONCE UPON A TIME...

→ Actually, in Pocono, 1948



Feynman & Schwinger present really strange methods to solve the infinities in QED

RENORMALIZATION

→ a set of arcane methods
to remove infinities

Example: Harmonic series

$$H(n) = 1 + \frac{1}{2} + \frac{1}{3} + \dots + \frac{1}{n}$$

$$\hookrightarrow \lim_{n \rightarrow \infty} H(n) \rightarrow \infty \quad (O(\log n))$$

↳ But if we use ζ -regularization

$$\zeta(s) = 1 + \frac{1}{2^s} + \frac{1}{3^s} + \dots$$

↳ Riemann zeta function

$$\lim_{n \rightarrow \infty} H(n) = \zeta(-1), \quad \zeta(-1) = -\frac{1}{12}$$

$$\Rightarrow \left[1 + \frac{1}{2} + \frac{1}{3} + \dots = -\frac{1}{12} \right]$$

Feynman described it as:

我们为求出 n 和 j 所玩的壳层游戏，在专业上叫做“重正化”（renormalization）。但是，不管这个词听来多聪明，我却说这个过程是蠢笨的！求助于这类戏法妨碍了我们去证明量子电动力学在数学上的自治性（self-consistent）。令人不解的是，尽管人们用了各种办法，这个理论至今仍未被证实是自治的；我猜想，重正化在数学上是不合法的。我们还没有一种好的数学方法描述量子电动力学，这是肯定的——像这样描述 n 、 j 同 m 、 e 之间关系的语言不是好的数学。 [23]

(From QED: The strange theory of light & matter / QED: 光和物质的奇妙理论)

- So it appears that the only things that depend on the small distances between coupling points are the values for n and j —theoretical numbers that are not directly observable any way; everything else, which can be observed, seems not to be affected. The shell game that we play to find n and j is technically called "renormalization." But no matter how clever the word, it is what I would call a dippy process! Having to resort to such hocus-pocus has prevented us from proving that the theory of quantum electrodynamics is mathematically self-consistent. It's surprising that the theory still hasn't been proved self-consistent one way or the other by now; I suspect that renormalization is not mathematically legitimate. What is certain is that we do not have a good mathematical way to describe the theory of quantum electrodynamics: such a bunch of words to describe the connection between n and j and m and e is not good mathematics.

◦ Richard Feynman, *QED: The Strange Theory of Light and Matter* (1985), Chap. 4. Loose Ends

↑ ORIGINAL FROM WIKIQUOTES

But after 70 years our understanding
of renormalization has vastly
improved

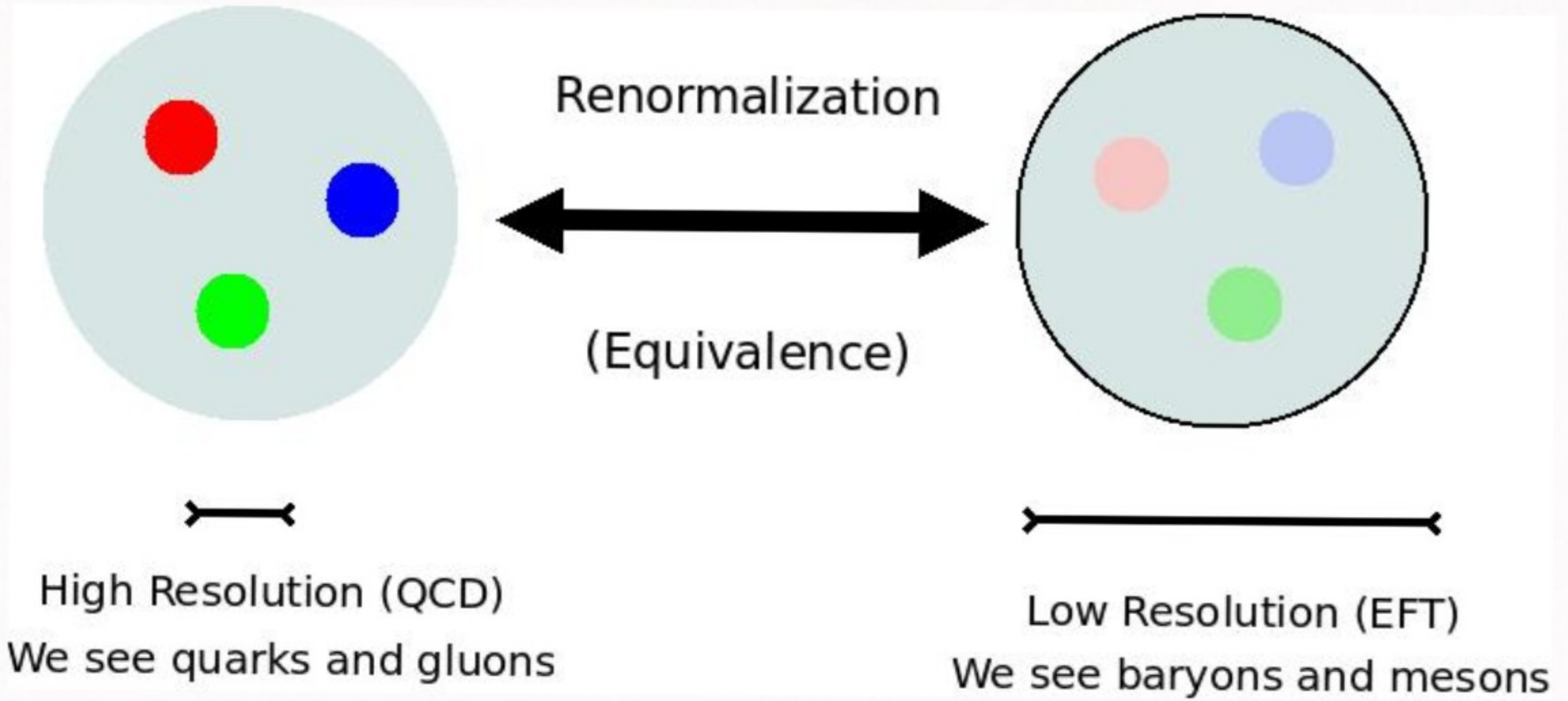


BASIC IDEA:

Physics at long-distances
does not depend on
short-distance details



Renormalization is just
the mathematical implementation
of this idea



↳ This is how it looks for hadron & nuclear physics

↳ RENORMALIZATION :

connects in a rigorous way the different views of a problem

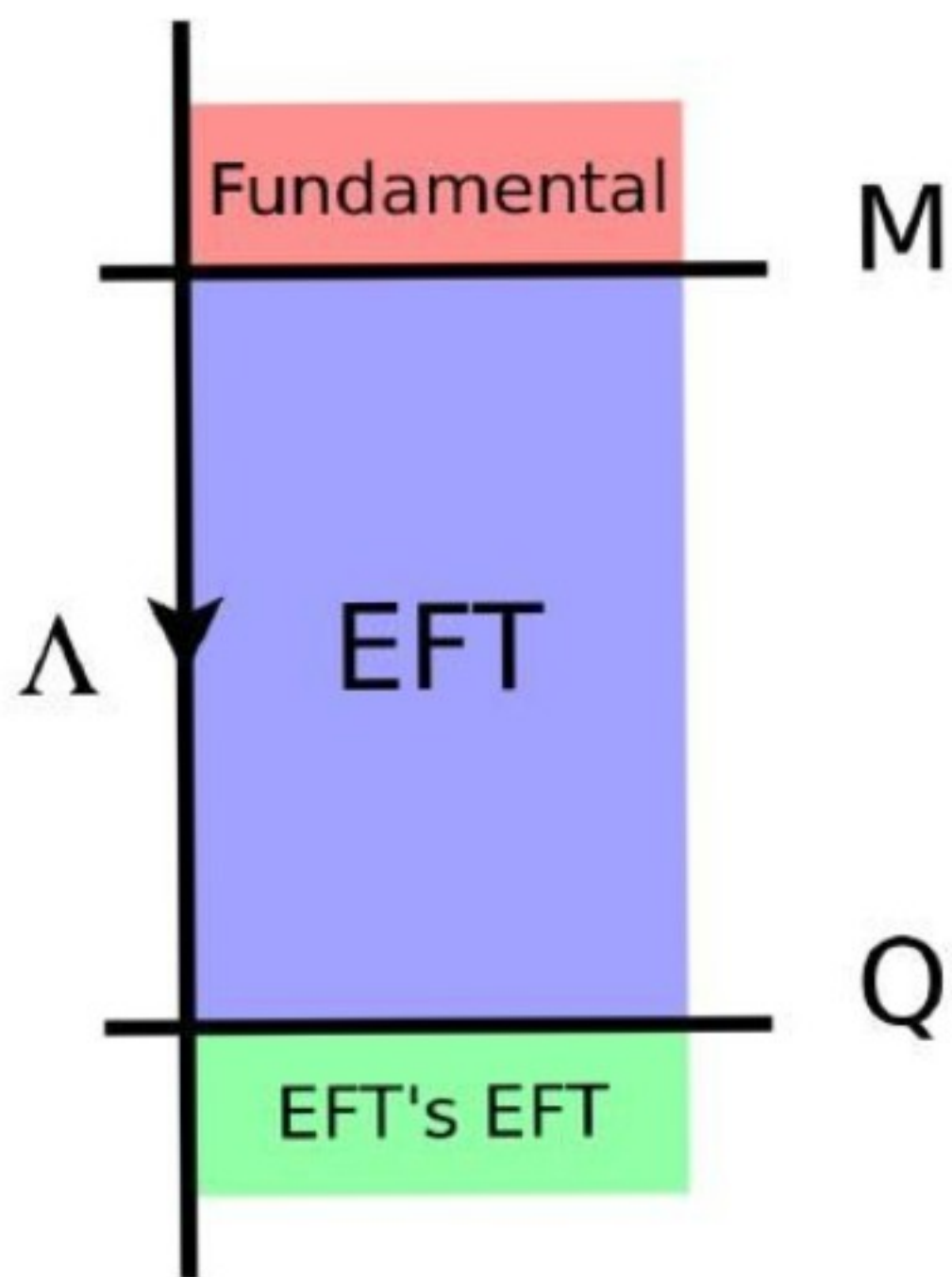
Quarks & gluons

$\Leftarrow \Rightarrow$

Baryons & mesons

RELATED CONCEPT:

Effective Field Theories (EFTs)



Physics is unique, but choice of theory depends on resolution Λ :

- ▶ $\Lambda \geq M$: Fundamental
- ▶ $M \geq \Lambda \geq Q$: EFT

For equivalent descriptions:

$$\frac{d}{d\Lambda} \langle \Psi | \mathcal{O} | \Psi \rangle = 0$$

Renormalization group invariance

[$\Lambda \rightarrow$ energy scale we are looking at
 $M \rightarrow$ scale of fundamental theory
 $Q \rightarrow$ scale of low energy theory
] THE CUTOFF

For QCD:

→ $M \sim 1 \text{ GeV}$ ($1/M \sim 0.2 \text{ fm}$)

mass of most hadrons

→ $Q \sim m_\pi \sim 140 \text{ MeV}$
($1/Q \sim 1.4 \text{ fm}$)

pion mass, average separation
of nucleons in nuclei



1) $\Lambda > M \rightarrow$ quark & gluon

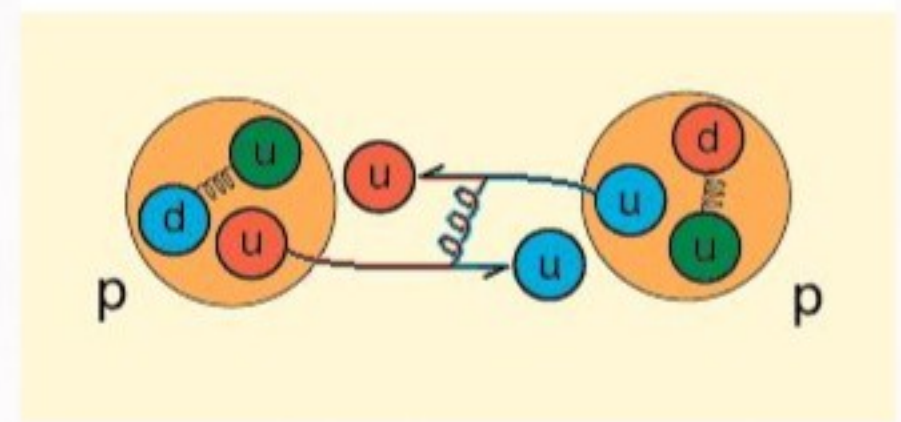
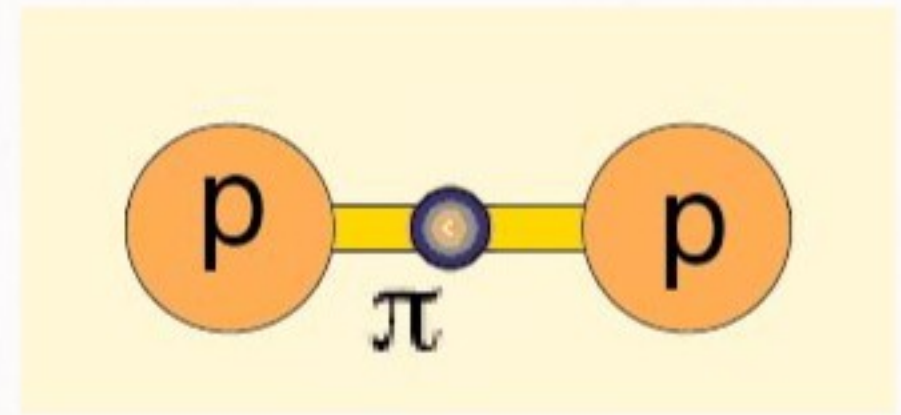
2) $\Lambda < M \rightarrow$ baryons
& meson

3) $\Lambda \sim M \rightarrow$ both descriptions
possible

This tells us how to build
an EFT:

1) $\Lambda \ll M$ two equivalent
descriptions:

baryons
& mesons



quarks &
gluons

2) If $\Lambda \rightarrow \infty$ both descriptions
will be equivalent if and only if
observables Λ -independent

$$\frac{d}{d\Lambda} \langle \Psi | \hat{O} | \Psi \rangle = 0$$

↑
wave function → Observable

But there is a problem

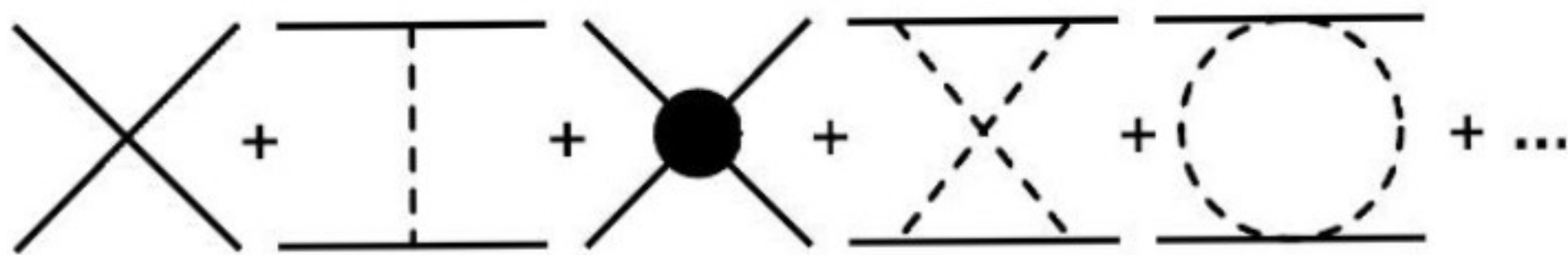
→ Predictive power

Begin at $\Lambda = M$, two equivalent descriptions

$$\underbrace{\text{quarks \& gluons}}_{\text{high energy}} \iff \underbrace{\text{hadrons (N, } \Delta, \pi)}_{\text{low energy}}$$

The hadron description equivalent if and only if

- (1) Include low energy symmetries (particularly **chiral symmetry**)
- (2) Consider infinite set of Feynman diagrams consistent with (1)



Problem: **infinite diagrams imply no predictive power**

1) FUNDAMENTAL DESCRIPTION

→ Quark & Gluons → QCD

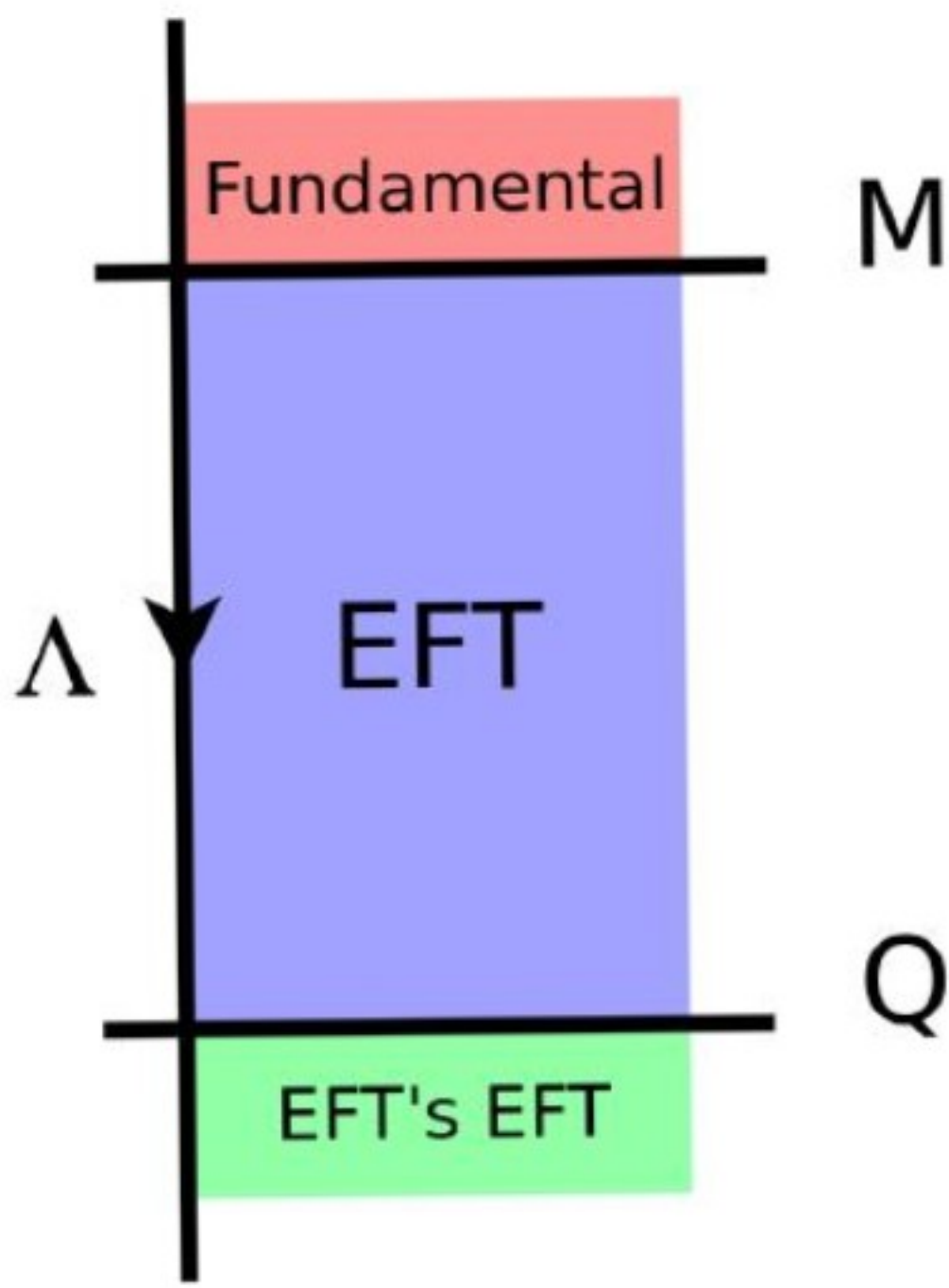
2) EFFECTIVE DESCRIPTION

→ Nucleons & pions

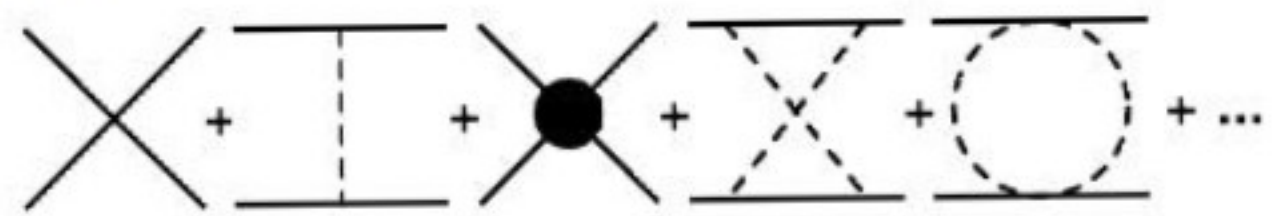
→ Only equivalent if we include all possible interactions (diagrams)

RECOVERING PREDICTIVE POWER

→ POWER COUNTING

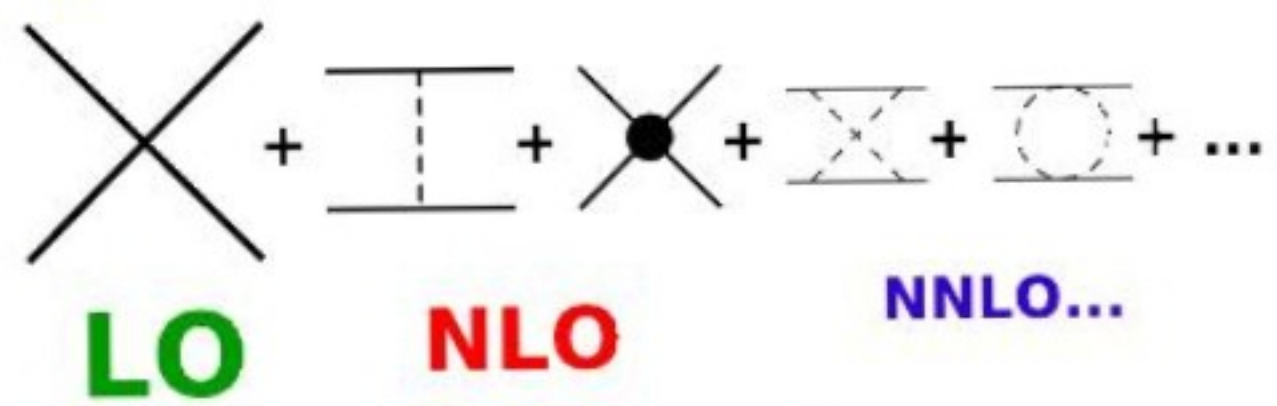


(1) At $\Lambda \sim M$ there is no order



$$\frac{d}{d\Lambda} \langle \Psi | \mathcal{O} | \Psi \rangle = 0$$

(2) while at $\Lambda \sim Q$ there is order



RGA implies that there's order at low energies / large distances



Power counting → some factors are more important than others

→ a bit too abstract,
so let's use an example

THE WONDERFUL THEORY OF TEACUPS & TEAPOTS



The good life TM

↳ a really good example
of an effective theory

A very innocent question :

Where does the tea cool's
faster ?



(1) The Teapot ?

(2) The Teacup ?

If your answer was (2):

Congratulations! You were right

↳ BUT WHY?



① TEAPOT :

a) Ceramic surface large

b) Exposed surface small

② TEACUP : a) small

b) large

HEAT TRANSFER :

[intuitively convection faster
than conduction]



[TEACUP COOLS FASTER
THAN TEAPOT]



But we didn't solve any equation
describing HEAT TRANSFER



Fundamental
theory



Effective
theory

FUNDAMENTAL THEORY:

1) Fourier's Law of heat conduction

$$\underline{q} = -k \nabla T$$

2) Convection-diffusion eq:

$$\frac{\partial c}{\partial t} = \nabla \cdot (D \nabla c) - \nabla \cdot (vc) + R$$

(taken from Wikipedia)

3) Use your computer and solve it for teapots & teacups



STRAIGHTFORWARD

BUT...

[I'm a really
lazy guy]

EFFECTIVE THEORY \rightarrow The smart way

1) Newton's law of cooling

$$T = T_0 e^{-\lambda(t-t_0)}$$

$T_0, t_0 \rightarrow$ initial temperature
of time

$T, t \rightarrow$ final temperature
of time

\rightarrow so I want to compute $\textcircled{1}$

2) Find the relevant degrees
of freedom



Exposed surface \textcircled{S}

Ceramic surface

$\textcircled{\Sigma}$

3) Propose a power counting

S more important than Σ

\Rightarrow expand in powers of

$$x = \left(\frac{S}{f\Sigma} \right) \quad \begin{array}{l} f: \text{correction} \\ \text{factor bc} \\ \text{ceramic coils} \\ S \text{ lower} \end{array}$$

4) Write down the theory

$$\lambda = S \left(c_0 + c_1 x + c_2 x^2 + \dots \right)$$

Low
Energy
Constants
(LECs)

it's a power series!
(that's why we call
it power counting)

5) Choose the accuracy we want Σ at the series

$$\mathcal{O}(x^0) : \lambda = c_0 \Sigma$$

↳ also called "Leading Order" (LO)

$$\mathcal{O}(x^1) : \lambda = c_0 \Sigma + c_1 \Sigma x$$

↳ also called "Next to Leading Order" (NLO)

$$\mathcal{O}(x^2) : \lambda = \Sigma (c_0 + c_1 x + c_2 x^2)$$

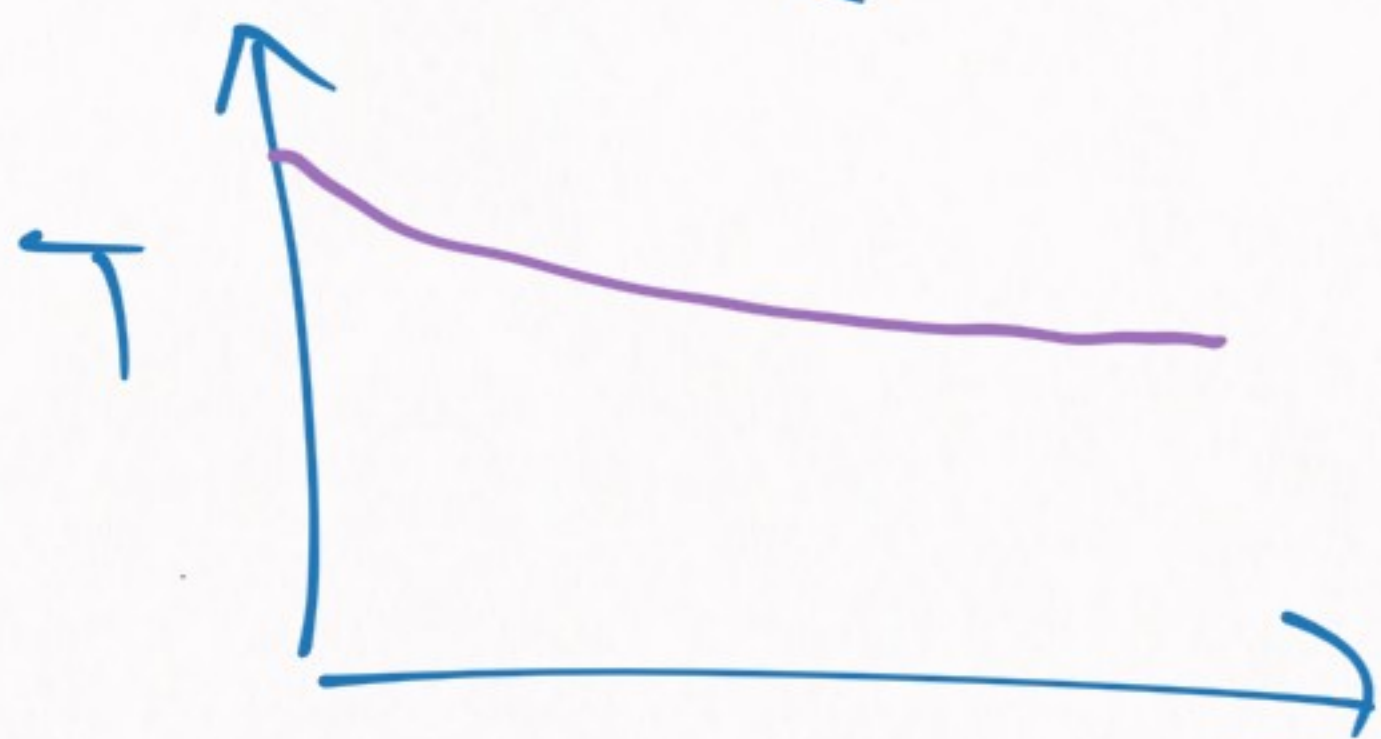
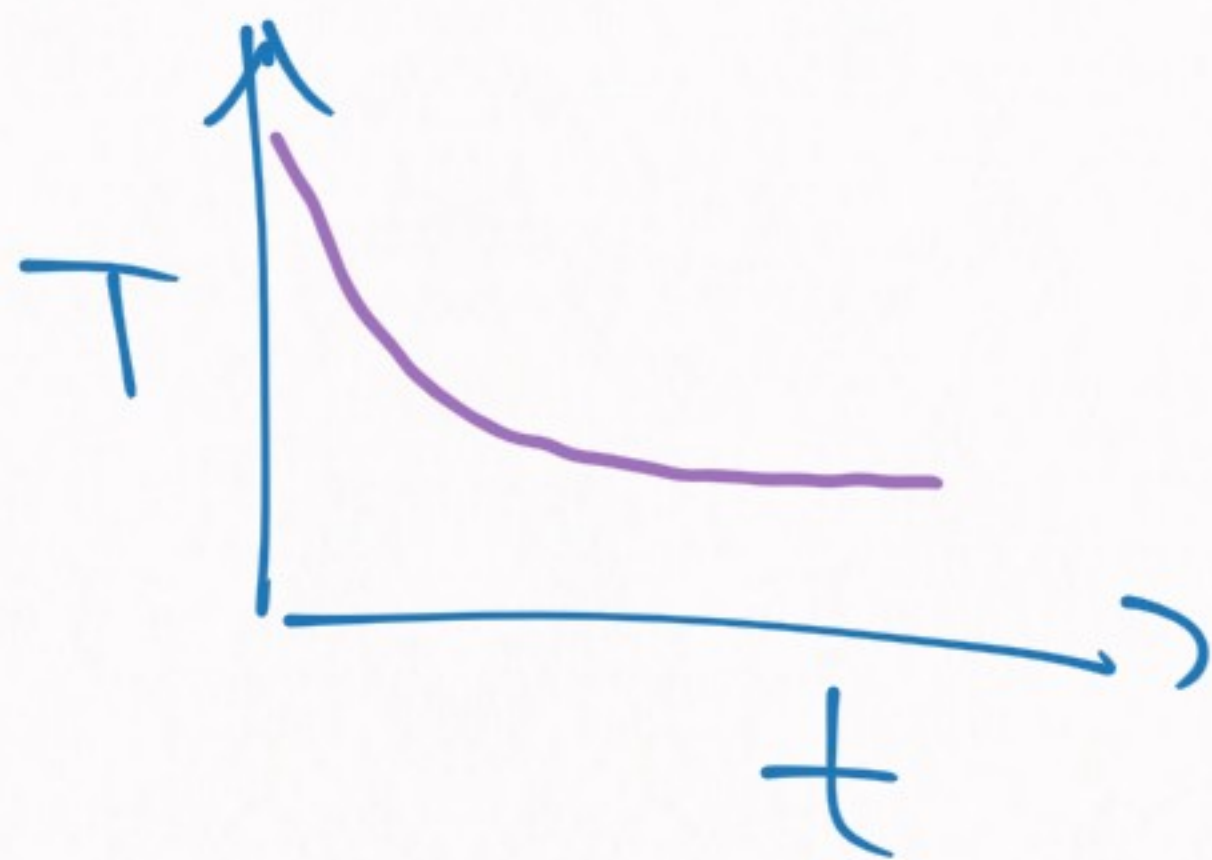
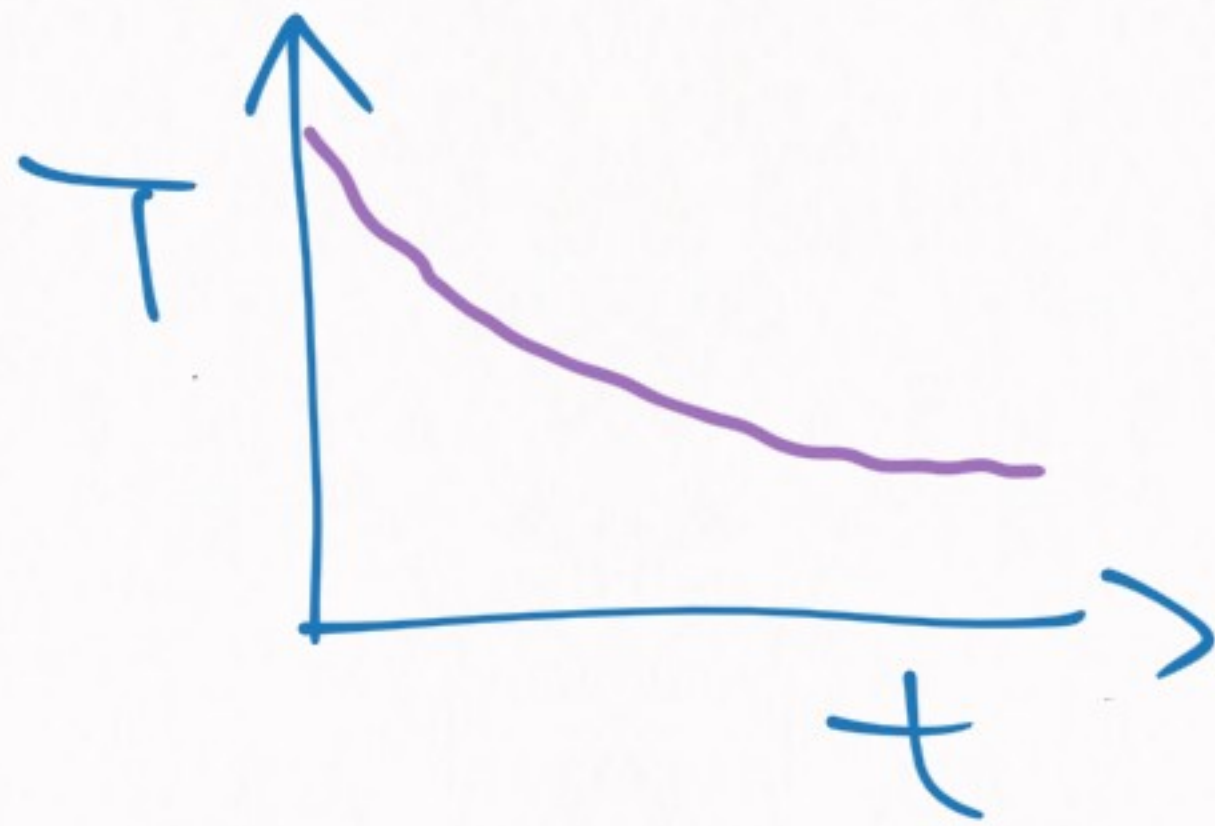
↳ NNLO or N²LO

More terms \Rightarrow more accuracy
(but also more difficult)

6) Fit the LECs (c_0, c_1, c_2, \dots)
to experimental data



→ my instruments



⇒ Fit the LECs ^{t}

RECAP:

- 1) Some dynamics (cooling's law
/ QM / AFT)
- 2) Degrees of freedom
(surfaces' types / particles / ...)
- 3) Power counting
- 4) Write down the theory
↳ LECs
- 5) Choose accuracy
↳ LO, NLO, N²LO, ...
- 6) Fit to experimental data

— ⊗ —

RELATIVELY SIMPLE

BUT THERE IS MORE

→ ∃ many power countings

Example: the teapot theory



↳ Σ

$$\lambda_{\text{tot}} = \sum \left(p_0 + \frac{p_1}{x} + \frac{p_2}{x^2} + \dots \right)$$
$$= \sum (p_0 + p_1 \gamma + p_2 \gamma^2 + \dots)$$

$$\gamma = \frac{p \Sigma}{S}, \quad \gamma < 1$$

Power counting not unique:

"Cup" power counting



$$\lambda_c = \sum (c_0 + c_1 x + c_2 x^2 + \dots)$$

(valid for $x < 1$)



"Pot" power counting

$$\lambda_p = \sum (p_0 + \frac{p_1}{x} + \frac{p_2}{x^2} + \dots)$$

(valid for $x > 1$)

In RG language we call these:



UNIVERSALTY CLASSES

OR

INFRARED FIXED POINTS

↳ they are the infrared ($\Lambda \rightarrow 0$) solutions of the RG (Renormalization Group)

More about RG language:

$\Lambda \rightarrow 0$ INFRARED ($r \rightarrow \infty$)

$\Lambda \rightarrow \infty$ ULTRAVIOLET
($r \rightarrow 0$)



Congratulations!

Now you know how
to renormalize
your tea set!



⇒ RENORMALIZED!

NEXT → We can try to renormalize
something related
w/ nuclear physics



A few recommendations:

1) The paper that gave us the RG

2. The Renormalization group and the epsilon expansion

(2700) K.G. Wilson (Princeton, Inst. Advanced Study & Cornell U., LNS), John B. Kogut

Published in **Phys.Rept.** 12 (1974) 75-199

DOI: [10.1016/0370-1573\(74\)90023-4](https://doi.org/10.1016/0370-1573(74)90023-4)

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[ADS Abstract Service](#)

[Detailed record](#) - [Cited by 2700 records](#) 1000+

↳ requires solid state
physics background

2) A wonderful exposition:

4. Renormalization and Effective Lagrangians

Joseph Polchinski ([Harvard U.](#)). Apr 1983. 27 pp.

Published in **Nucl.Phys.** B231 (1984) 269-295

HUTP-83-A018

DOI: [10.1016/0550-3213\(84\)90287-6](https://doi.org/10.1016/0550-3213(84)90287-6)

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[KEK scanned document](#)

[Detailed record](#) - [Cited by 1120 records](#) 1000+

↳ requires QFT & d^4 theory

3) This one is awesome & easy

1. Building light nuclei from neutrons, protons, and pions

(45) Daniel R. Phillips (Ohio U.). Mar 2002. 54 pp.

Published in **Czech.J.Phys.** 52 (2002) B49

DOI: [10.1007/s10582-002-0079-z](https://doi.org/10.1007/s10582-002-0079-z)

To appear in the proceedings of Conference: [C01-07-09.13 Proceedings](#)

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4) A bit more difficult, but worth a try

3. A Renormalization group treatment of two-body scattering

(196) Michael C. Birse, Judith A. McGovern, Keith G. Richardson (Manchester U.). Jul 1998. 4 pp.

Published in **Phys.Lett.** B464 (1999) 169-176

MC-TH-98-11

DOI: [10.1016/S0370-2693\(99\)00991-0](https://doi.org/10.1016/S0370-2693(99)00991-0)

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[ADS Abstract Service](#)

[Detailed record](#) - [Cited by 196 records](#) 100+

5) A classic of nuclear EFT

1. Effective field theory for few nucleon systems

(618) Paulo F. Bedaque (LBL, Berkeley), Udirajara van Kolck (Arizona U. & RIKEN BNL). Mar 2002. 55 pp.

Published in **Ann.Rev.Nucl.Part.Sci.** 52 (2002) 339-396

DOI: [10.1146/annurev.nucl.52.050102.090637](https://doi.org/10.1146/annurev.nucl.52.050102.090637)

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[ADS Abstract Service](#)

[Detailed record](#) - [Cited by 618 records](#) 500+

the big boss

6) Even I have a short review

23. Power Counting and Wilsonian Renormalization in Nuclear Effective Field Theory

(28) Manuel Pavon Valderrama (Beihang U. & Orsay, IPN). Apr 5, 2016. 40 pp.

Published in **Int.J.Mod.Phys.** E25 (2016) no.05, 1641007

DOI: [10.1142/S021830131641007X](https://doi.org/10.1142/S021830131641007X)

e-Print: [arXiv:1604.01332](#) [nucl-th] | [PDF](#)

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